

Optics I: lenses and apertures

CS 178, Spring 2014

Begun 4/8/14, finished 4/10.

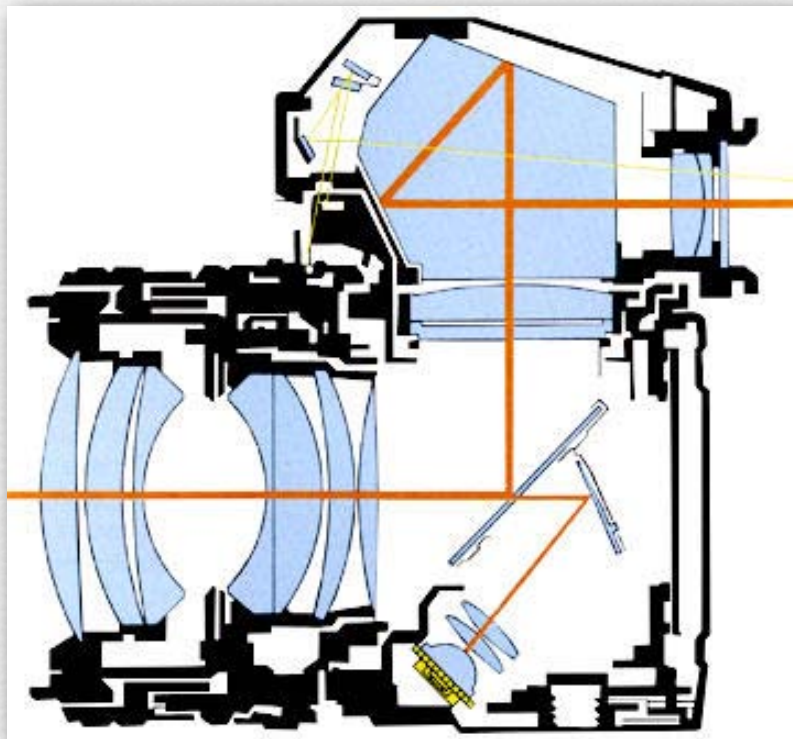


Marc Levoy
Computer Science Department
Stanford University

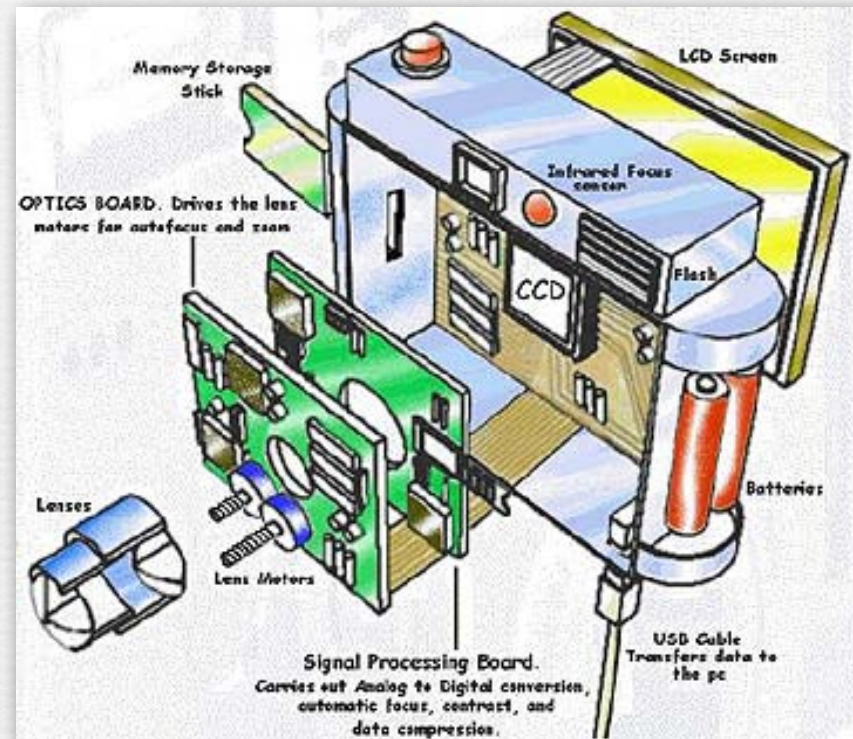
Outline

- ◆ why study lenses?
 - ◆ thin lenses
 - graphical constructions, algebraic formulae
 - ◆ thick lenses
 - center of perspective, lens as $3D \rightarrow 3D$ transformation
 - ◆ depth of field
-
- ◆ aberrations & distortion
 - ◆ vignetting, glare, and other lens artifacts
 - ◆ diffraction and lens quality
 - ◆ special lenses
 - telephoto, zoom

Cameras and their lenses

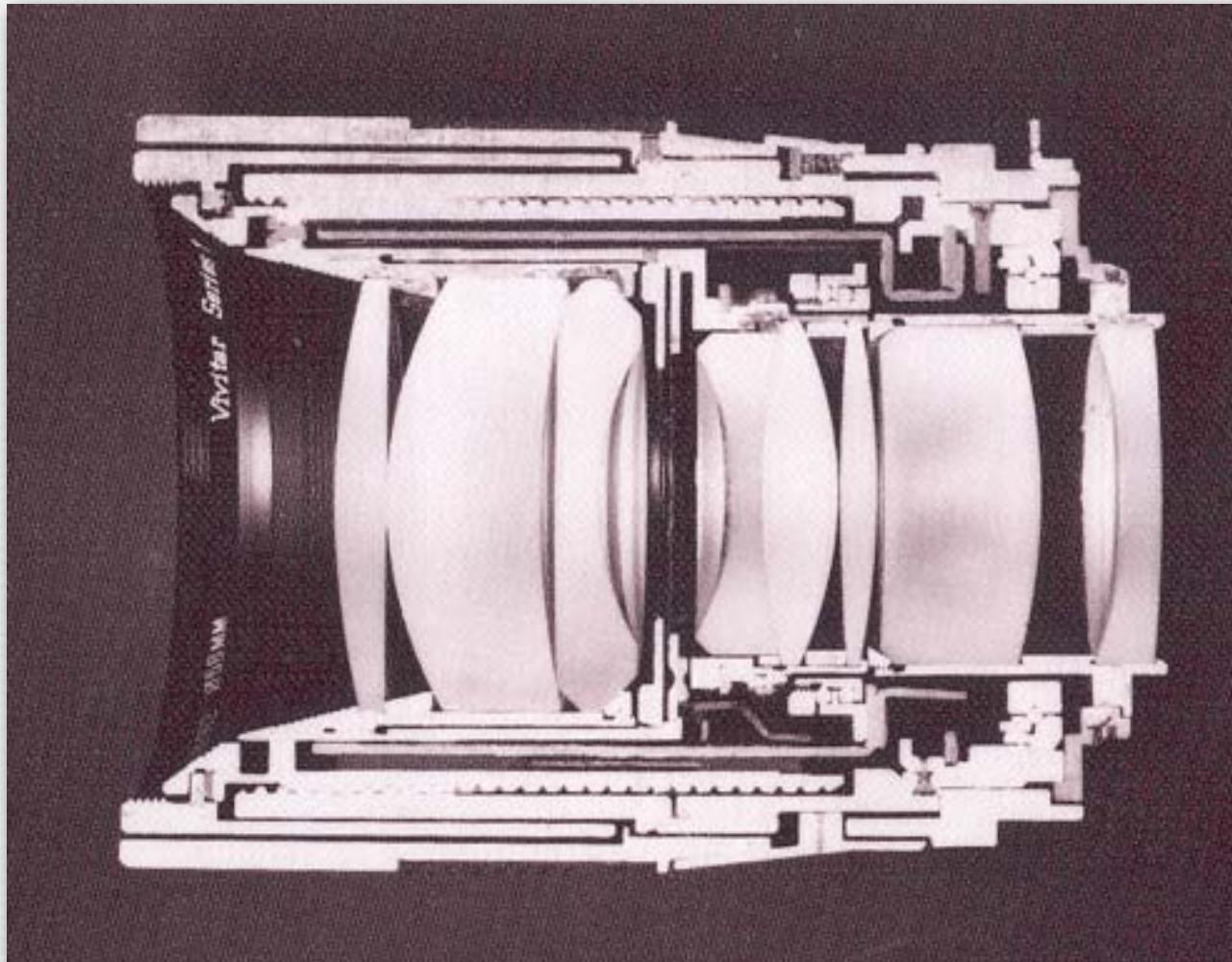


single lens reflex
(SLR) camera



digital still camera (DSC),
i.e. point-and-shoot

Cutaway view of a real lens



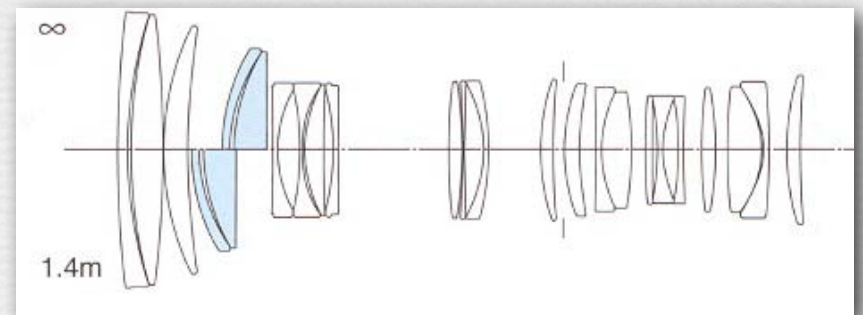
Vivitar Series 1 90mm f/2.5
Cover photo, Kingslake, *Optics in Photography*

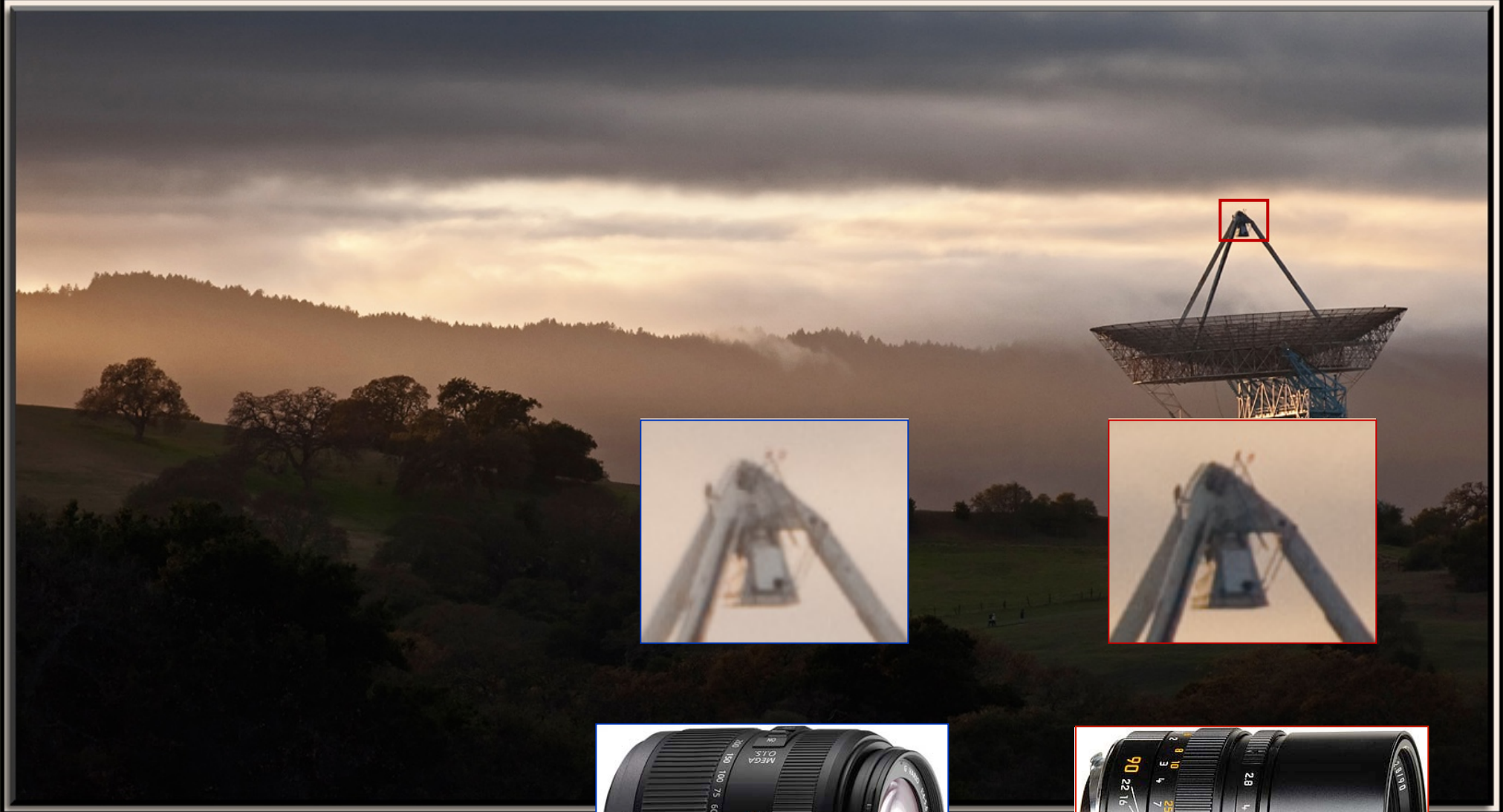
Lens quality varies

- ◆ Why is this toy so expensive?
 - EF 70-200mm f/2.8L IS USM
 - \$1700

- ◆ Why is it better than this toy?
 - EF 70-300mm f/4-5.6 IS USM
 - \$550

- ◆ And why is it so complicated?





Panasonic 45-200/4-5.6
zoom, at 200mm f/4.6
\$300



Leica 90mm/2.8 Elmarit-M
prime, at f/4
\$2000

Stanford Big Dish
Panasonic GF1

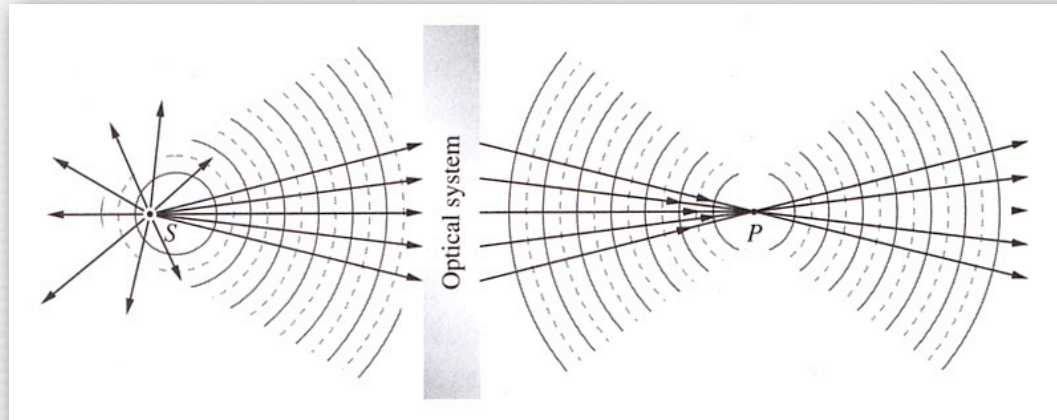
Zoom lens versus prime lens



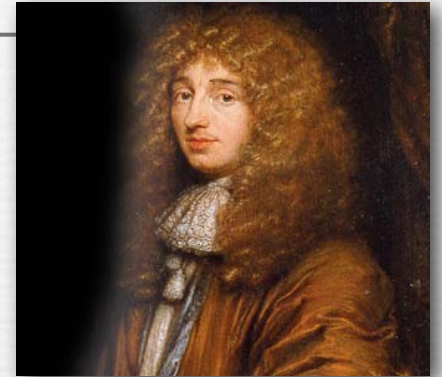
Canon 100-400mm/4.5-5.6
zoom, at 300mm and f/5.6
\$1600

Canon 300mm/2.8
prime, at f/5.6
\$4300

Physical versus geometrical optics



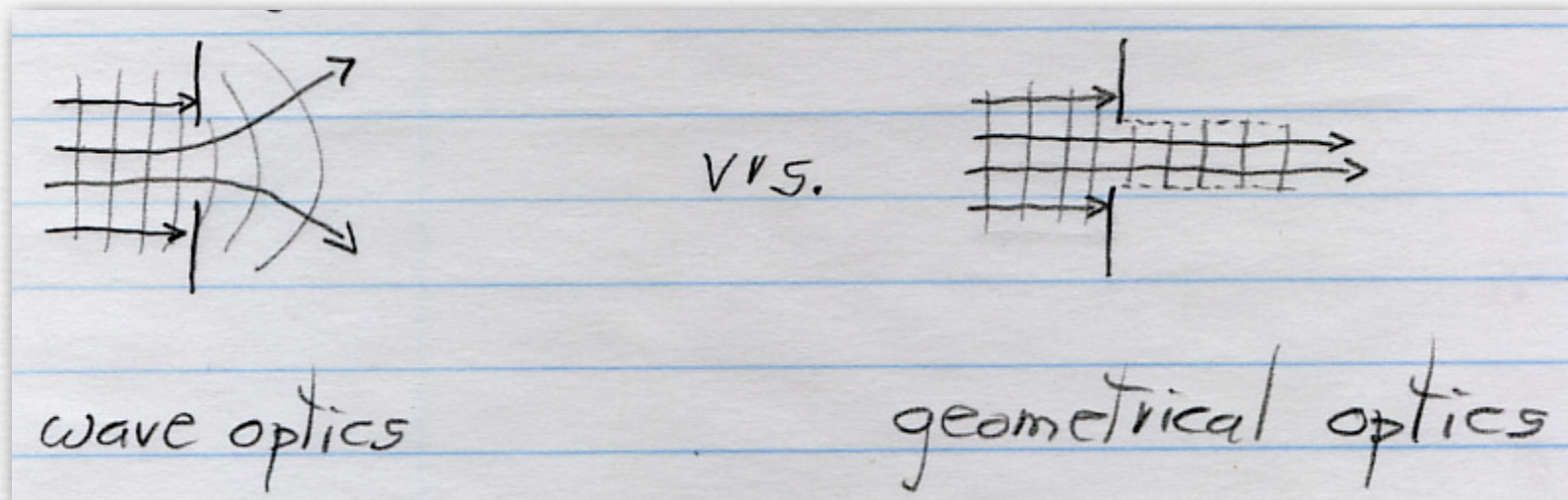
(Hecht)



Christiaan Huygens

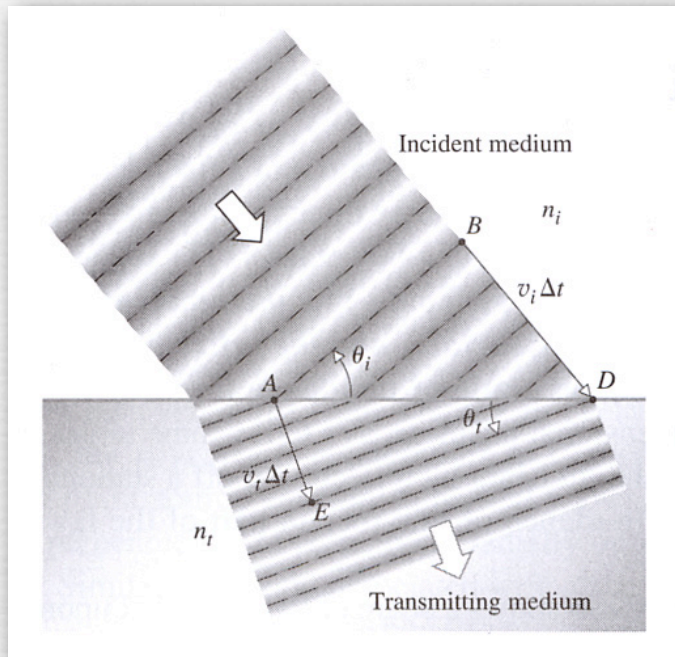
- ◆ light can be modeled as traveling waves
- ◆ the perpendiculars to these waves can be drawn as rays
- ◆ diffraction causes these rays to bend, e.g. at a slit
- ◆ *geometrical optics* assumes
 - $\lambda \rightarrow 0$
 - no diffraction
 - in free space, rays are straight (a.k.a. rectilinear propagation)

Physical versus geometrical optics (contents of whiteboard)

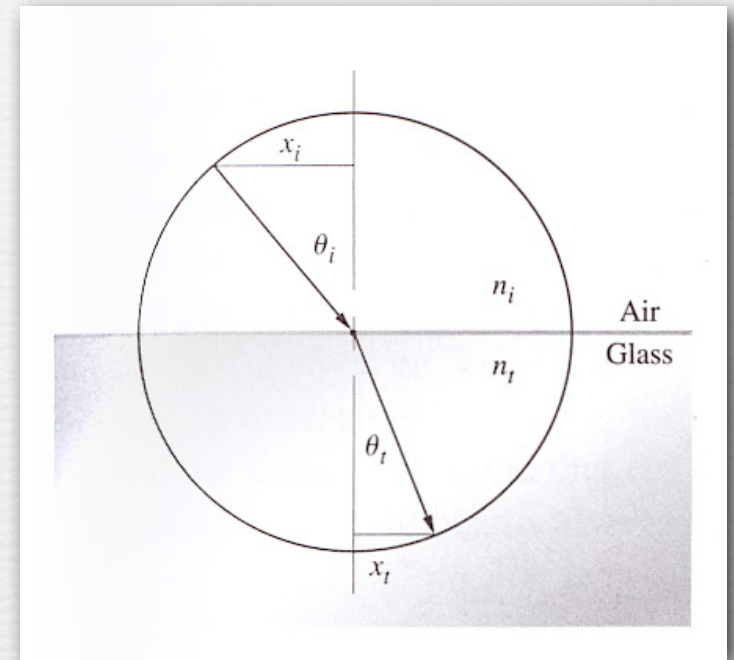


- ◆ in geometrical optics, we assume that rays do not bend as they pass through a narrow slit
- ◆ this assumption is valid if the slit is much larger than the wavelength, represented on the previous slide by the limit $\lambda \rightarrow 0$
- ◆ physical optics is a.k.a. wave optics

Snell's law of refraction



(Hecht)



- ◆ as waves change speed at an interface, they also change direction
- ◆ index of refraction n_r is defined as

$$\frac{x_i}{x_t} = \frac{\sin \theta_i}{\sin \theta_r} = \frac{n_t}{n_i}$$

$\frac{\text{speed of light in a vacuum}}{\text{speed of light in medium } r}$

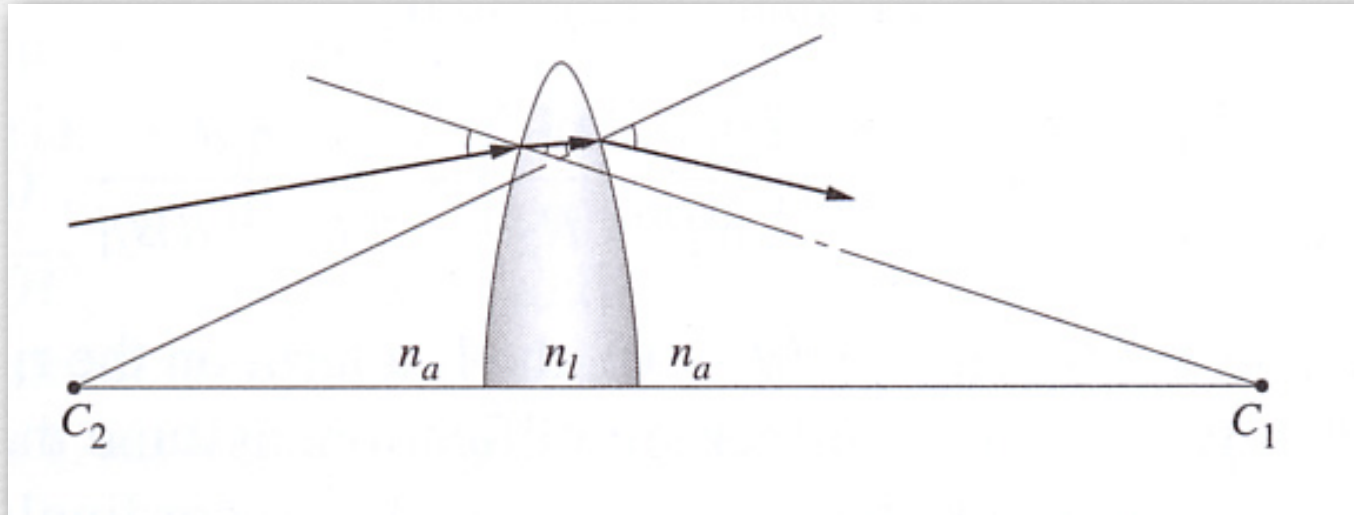
Typical refractive indices (n)

- ◆ air = ~ 1.0
- ◆ water = 1.33
- ◆ glass = 1.5 - 1.8



mirage due to changes in the index of refraction of air with temperature

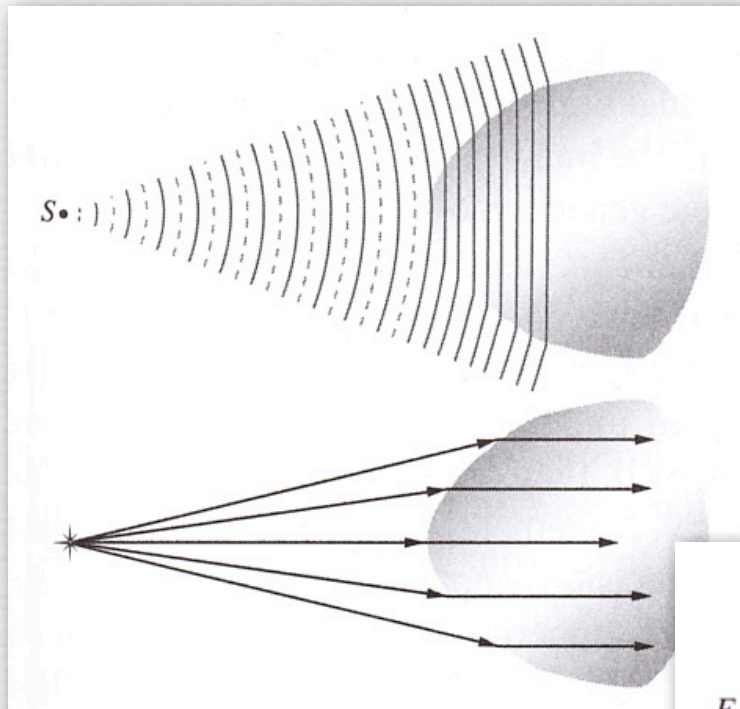
Refraction in glass lenses



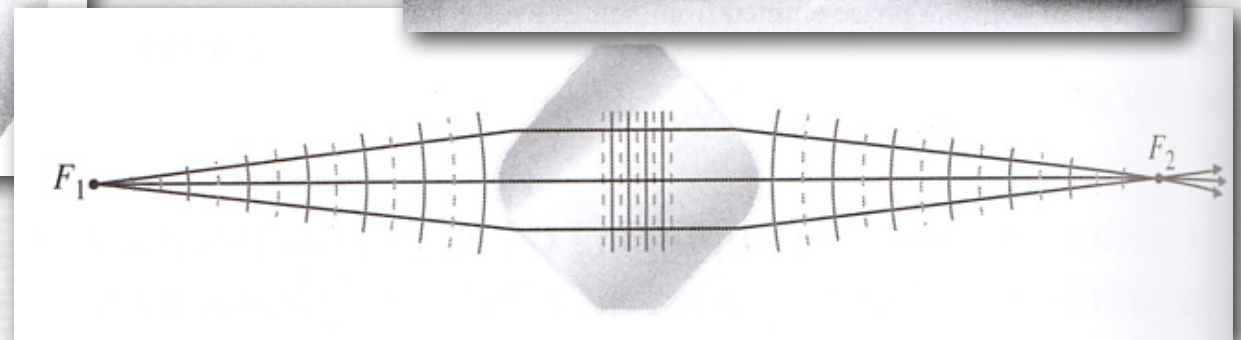
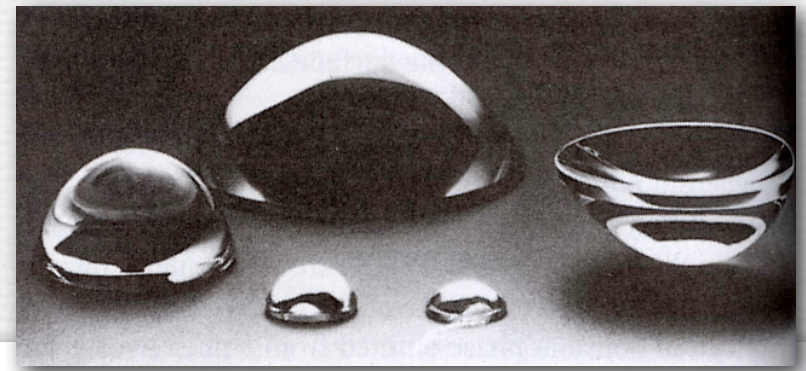
(Hecht)

- ◆ when transiting from air to glass, light bends towards the normal
- ◆ when transiting from glass to air, light bends away from the normal
- ◆ light striking a surface perpendicularly does not bend

Q. What shape should a refractive interface be to make parallel rays converge to a point?



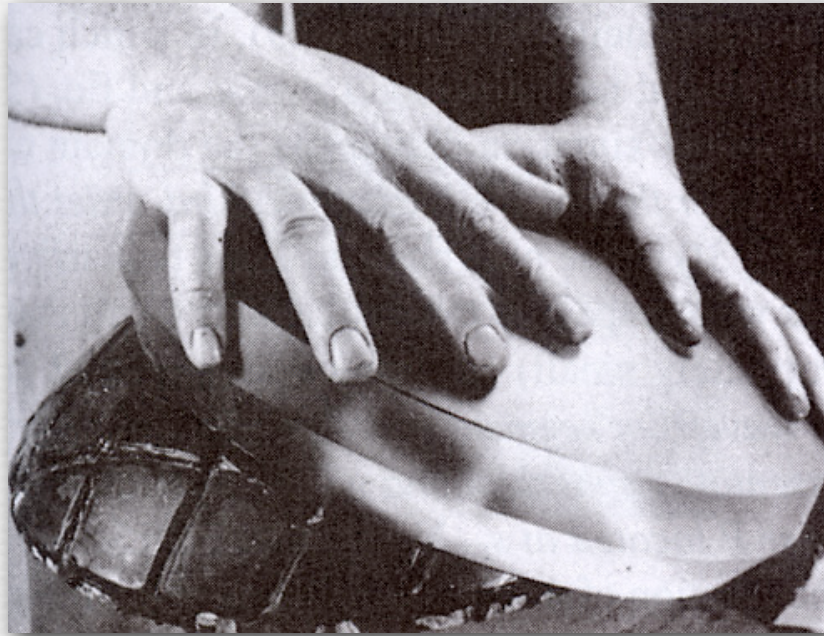
(Hecht)



A. a hyperbola

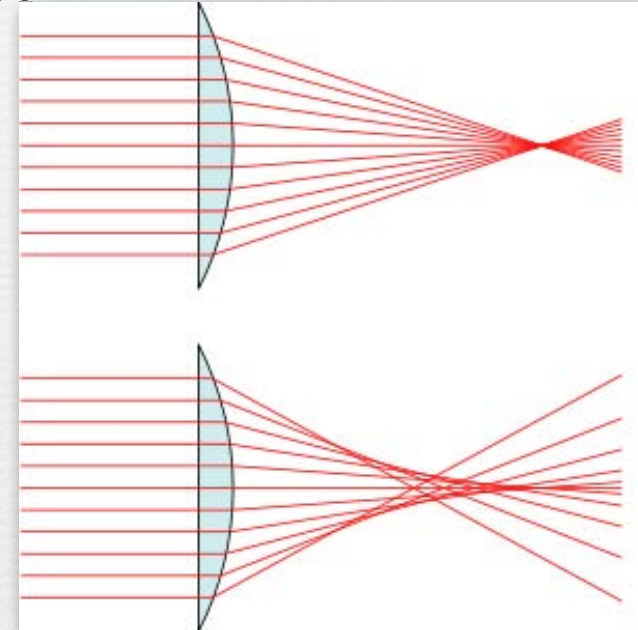
◆ so lenses should be hyperbolic!

Spherical lenses



(Hecht)

hyperbolic lens



spherical lens

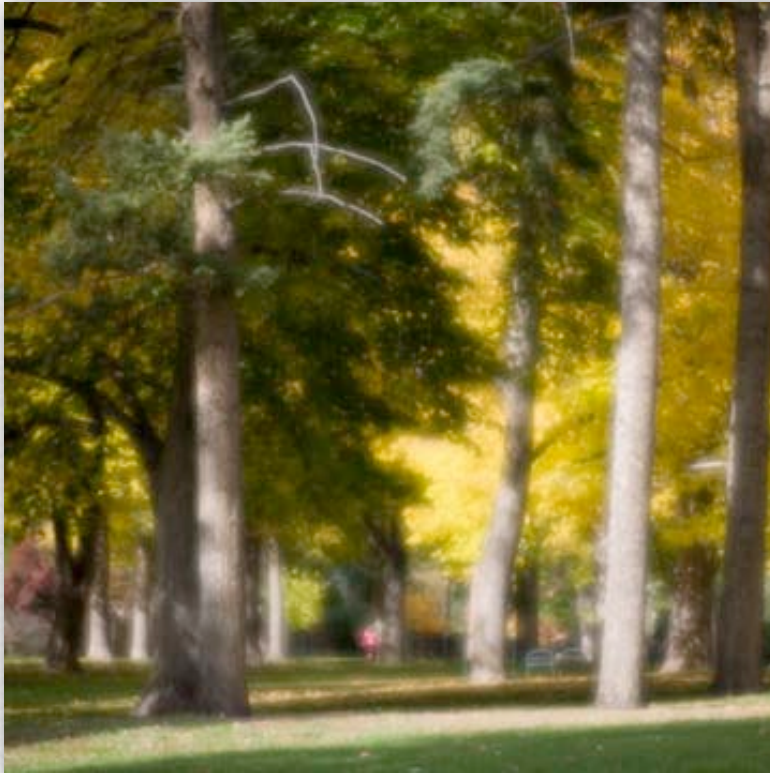
(wikipedia)

- ◆ two roughly fitting curved surfaces ground together will eventually become spherical
- ◆ spheres don't bring parallel rays to a point
 - this is called *spherical aberration*
 - nearly axial rays (*paraxial rays*) behave best

Examples of spherical aberration



Canon 135mm
soft focus lens



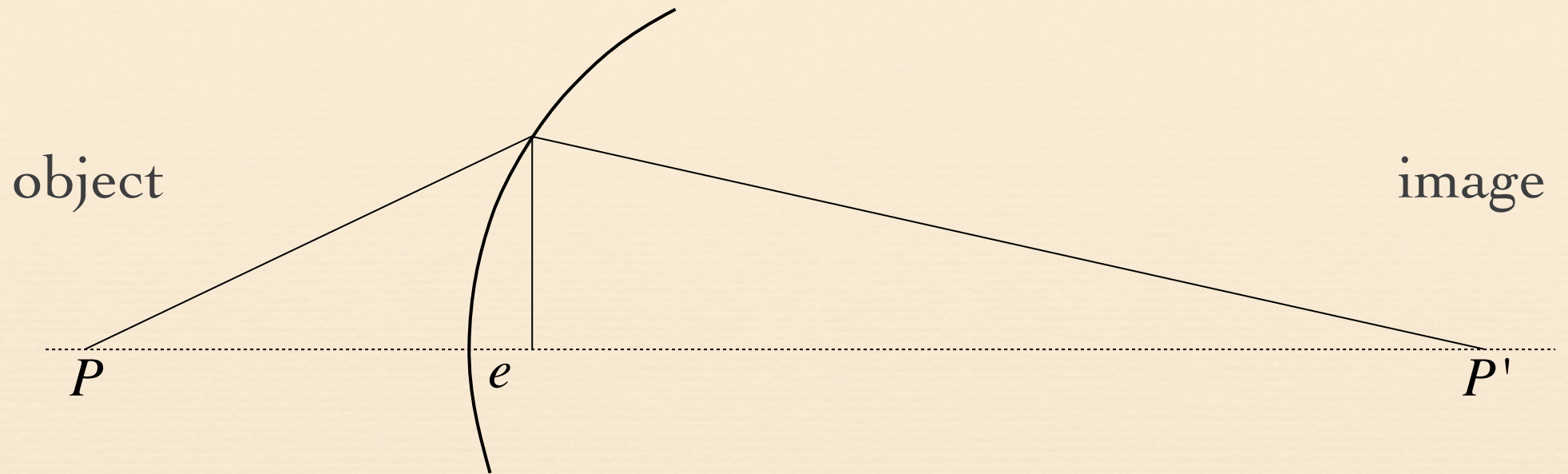
(gtmerideth)



(Canon)



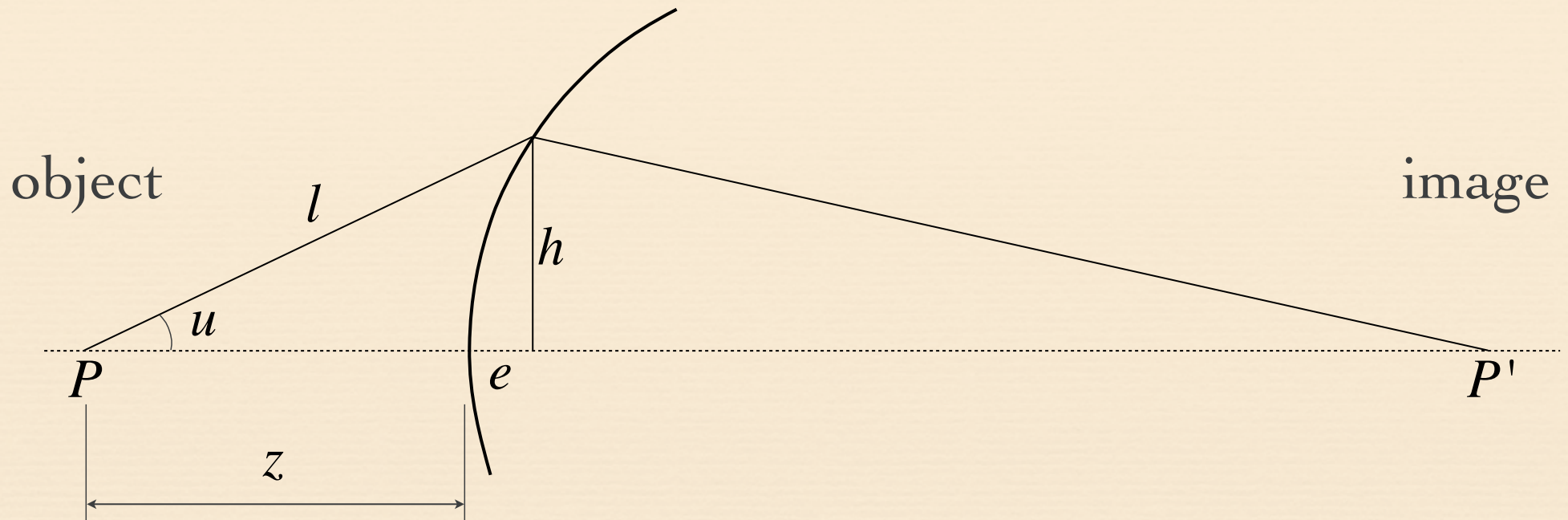
Paraxial approximation



♦ assume $e \approx 0$

Not responsible on exams
for orange-tinted slides

Paraxial approximation



- ◆ assume $e \approx 0$
- ◆ assume $\sin u = h/l \approx u$ (for u in radians)
- ◆ assume $\cos u \approx z/l \approx 1$
- ◆ assume $\tan u \approx \sin u \approx u$

The paraxial approximation is a.k.a. first-order optics

◆ assume first term of $\sin \phi = \phi - \frac{\phi^3}{3!} + \frac{\phi^5}{5!} - \frac{\phi^7}{7!} + \dots$

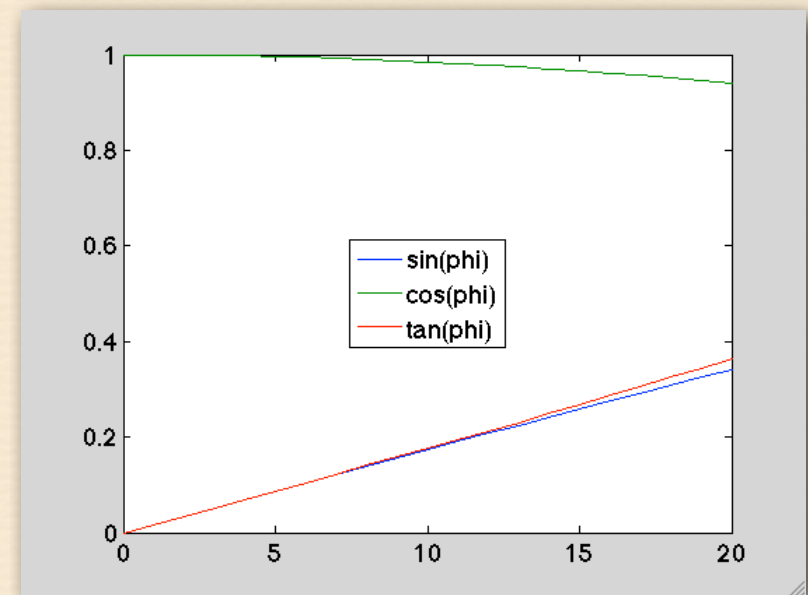
- i.e. $\sin \phi \approx \phi$

◆ assume first term of $\cos \phi = 1 - \frac{\phi^2}{2!} + \frac{\phi^4}{4!} - \frac{\phi^6}{6!} + \dots$

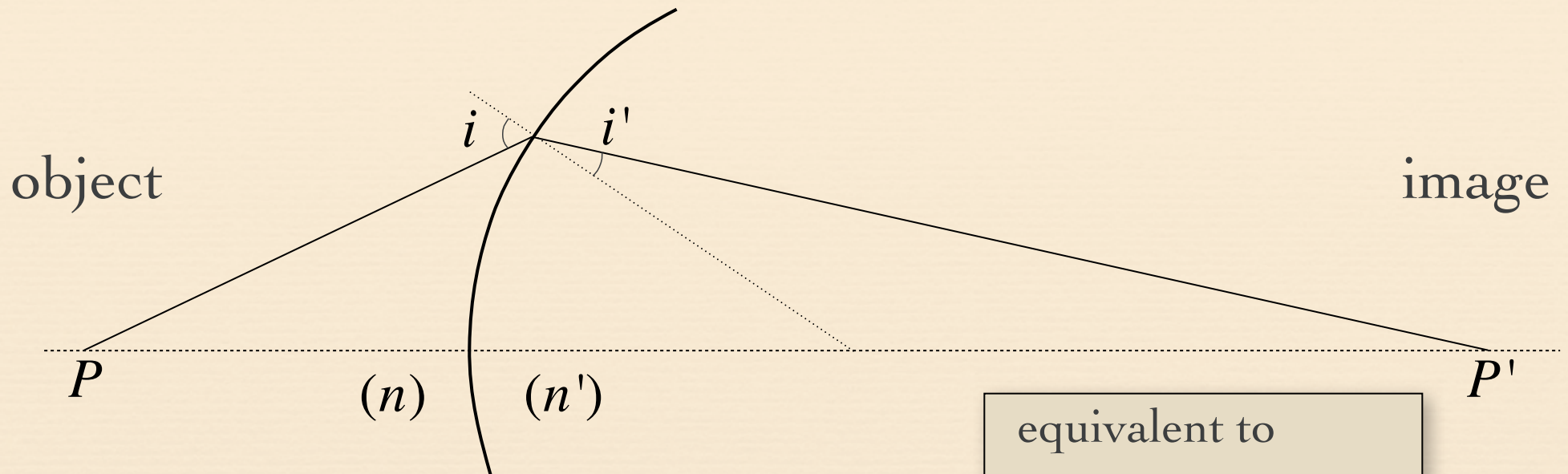
- i.e. $\cos \phi \approx 1$
- so $\tan \phi \approx \sin \phi \approx \phi$

these are the
Taylor series for
 $\sin \phi$ and $\cos \phi$

(phi in degrees)



Paraxial focusing



Snell's law:

$$n \sin i = n' \sin i'$$

paraxial approximation:

$$n i \approx n' i'$$

equivalent to

$$\frac{\sin \theta_i}{\sin \theta_t} = \frac{n_t}{n_i}$$

with

$$n = n_i \text{ for air}$$

$$n' = n_t \text{ for glass}$$

i, i' in radians

θ_i, θ_t in degrees

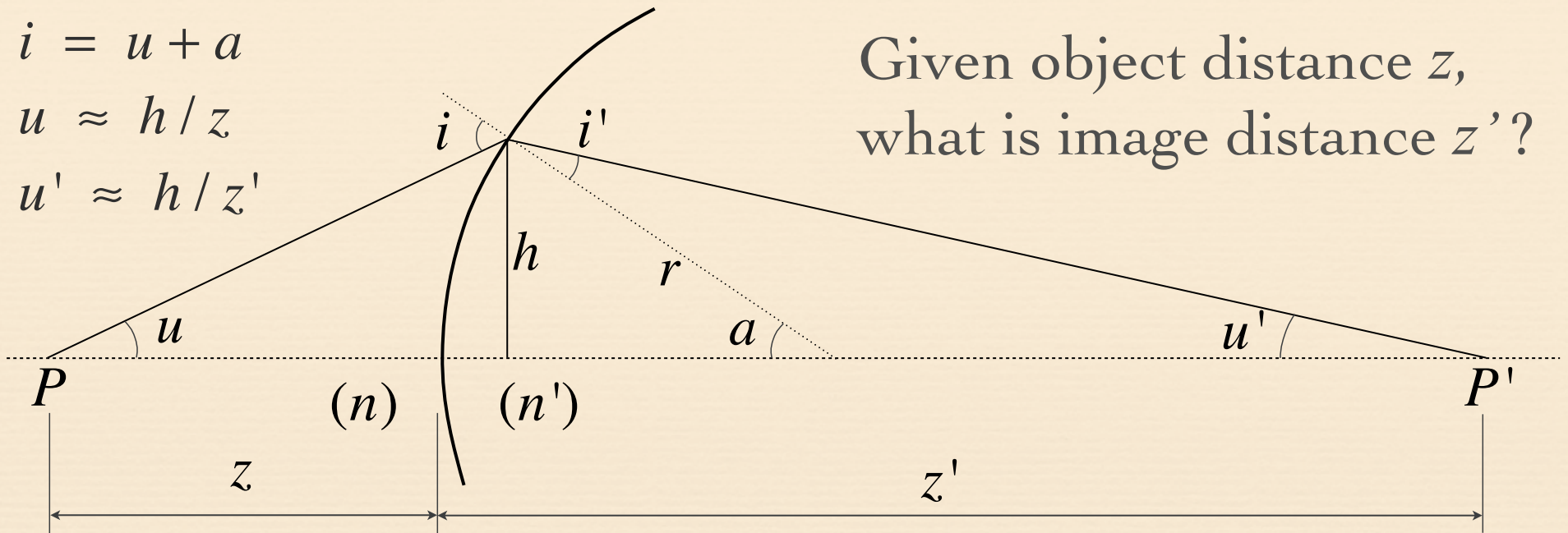
Paraxial focusing

$$i = u + a$$

$$u \approx h/z$$

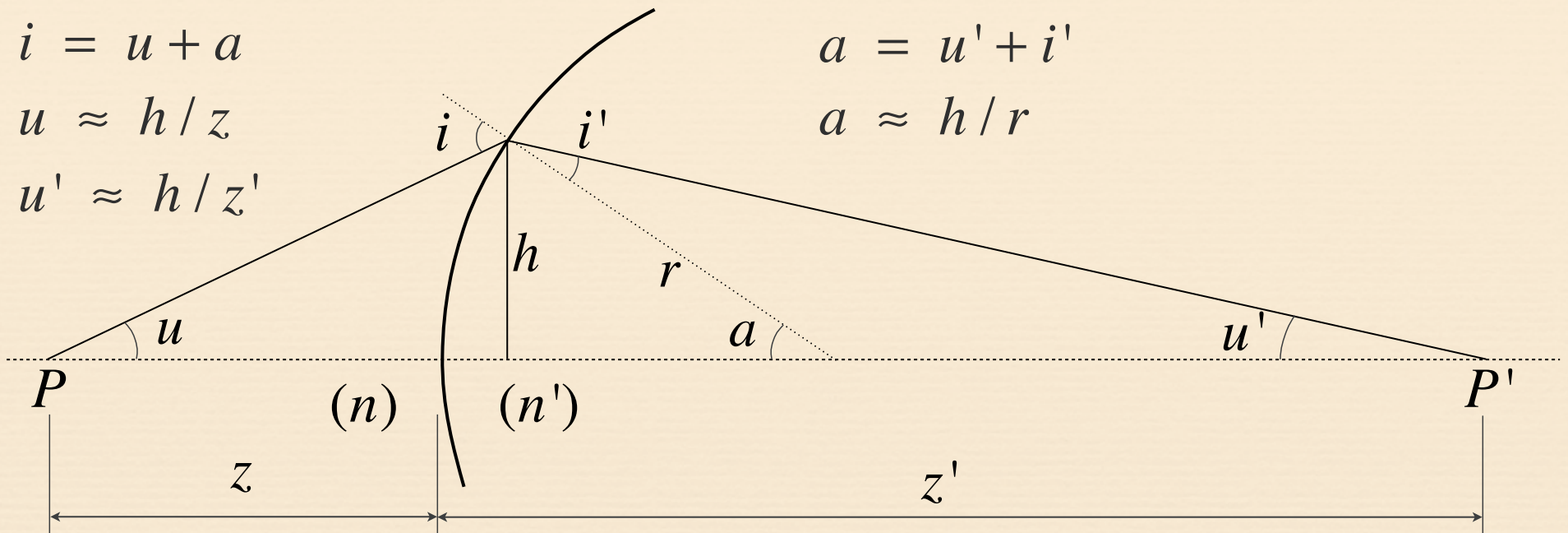
$$u' \approx h/z'$$

Given object distance z ,
what is image distance z' ?



$$n i \approx n' i'$$

Paraxial focusing



$$n(u + a) \approx n'(a - u')$$

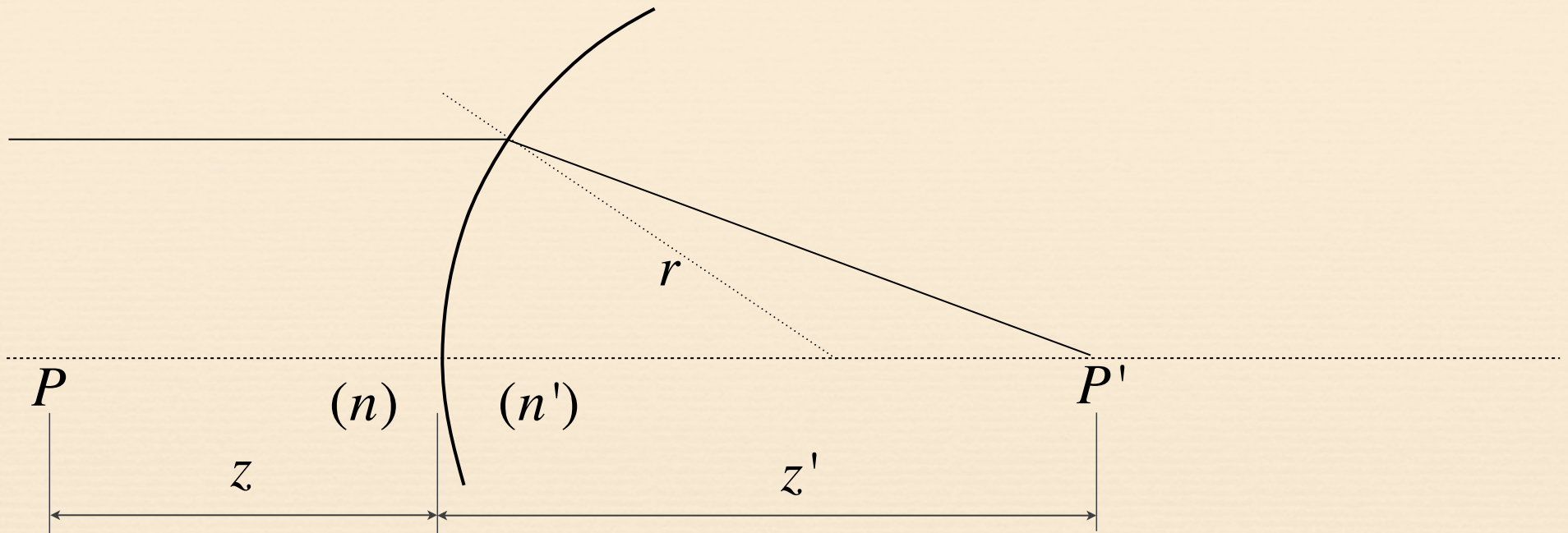
$$n(h/z + h/r) \approx n'(h/r - h/z')$$

$$n/z + n/r \approx n'/r - n'/z'$$

$$ni \approx n'i'$$

◆ h has canceled out, so any ray from P will focus to P'

Focal length



What happens if z is ∞ ?

$$n/z + n/r \approx n'/r - n'/z'$$

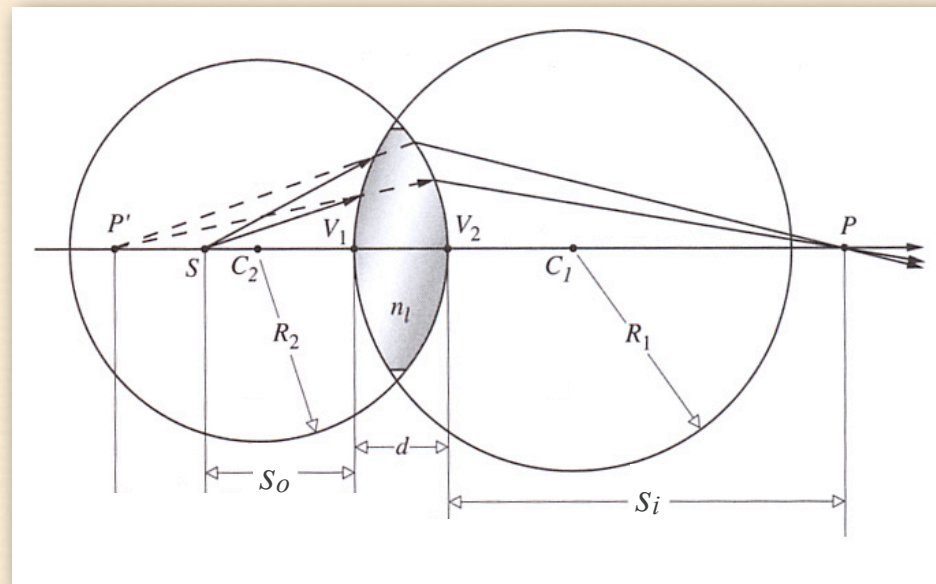
$$n/r \approx n'/r - n'/z'$$

$$z' \approx (r n') / (n' - n)$$

♦ $f \triangleq$ focal length = z'

Lensmaker's formula

- ♦ using similar derivations, one can extend these results to two spherical interfaces forming a lens in air



(Hecht, edited)

- ♦ as $d \rightarrow 0$ (*thin lens approximation*), we obtain the lensmaker's formula

$$\frac{1}{s_o} + \frac{1}{s_i} = (n_l - 1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right)$$

Gaussian lens formula

- ◆ Starting from the lensmaker's formula

$$\frac{1}{s_o} + \frac{1}{s_i} = (n_l - 1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right), \quad (\text{Hecht, eqn 5.15})$$

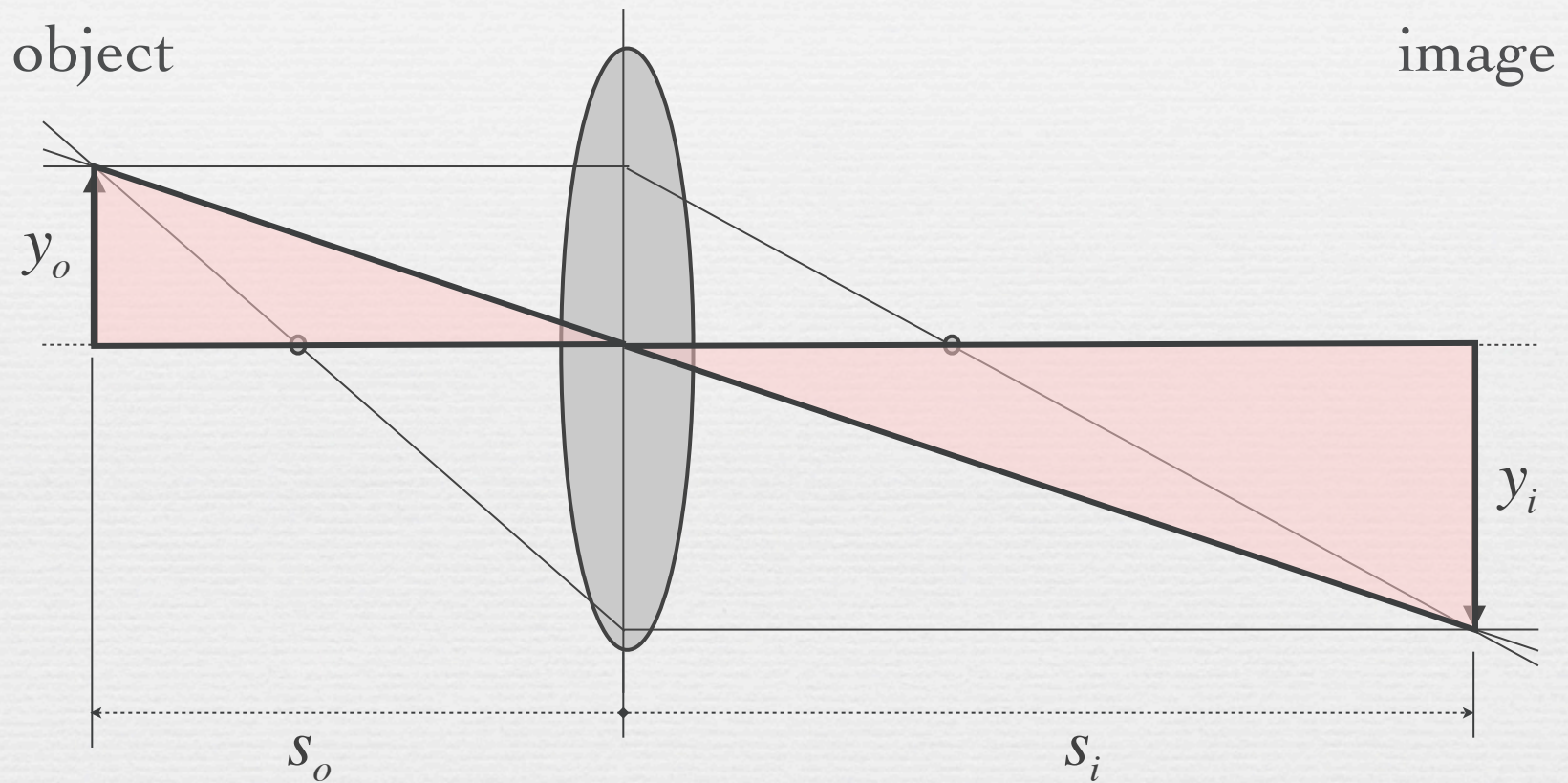
- ◆ and recalling that as object distance s_o is moved to infinity, image distance s_i becomes focal length f_i , we get

$$\frac{1}{f_i} = (n_l - 1) \left(\frac{1}{R_1} - \frac{1}{R_2} \right). \quad (\text{Hecht, eqn 5.16})$$

- ◆ Equating these two, we get the Gaussian lens formula

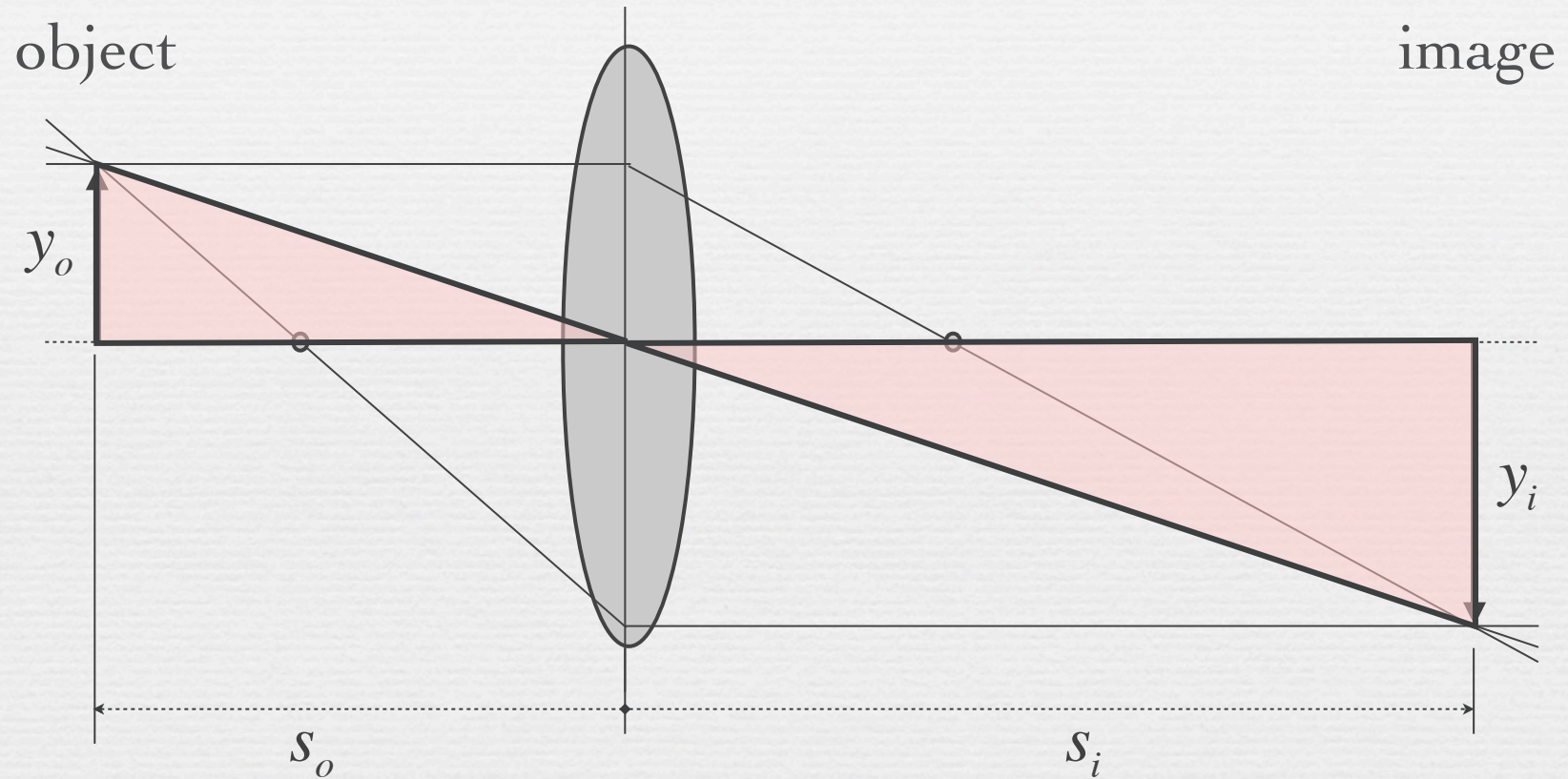
$$\frac{1}{s_o} + \frac{1}{s_i} = \frac{1}{f_i}. \quad (\text{Hecht, eqn 5.17})$$

From Gauss's ray construction to the Gaussian lens formula



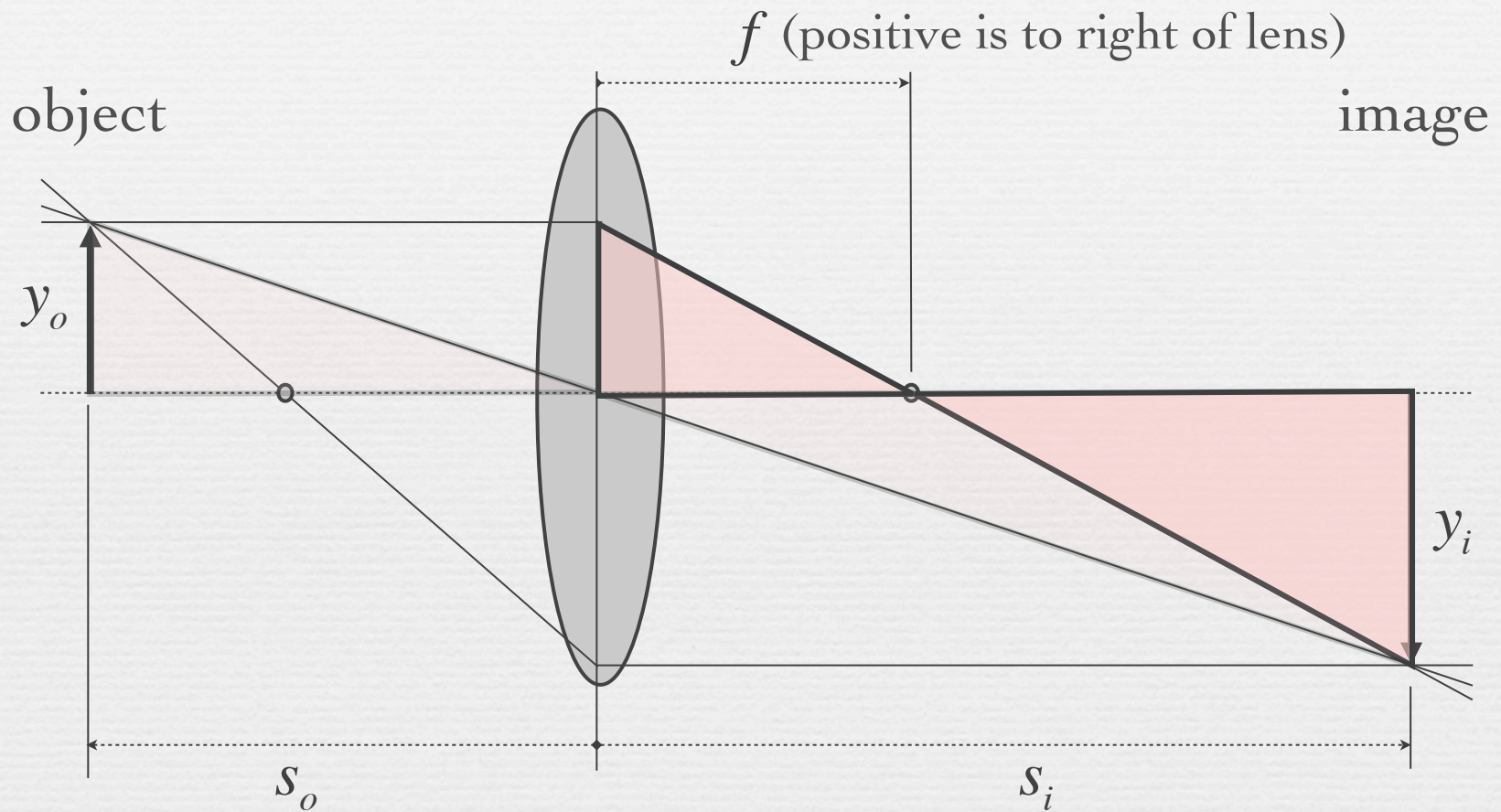
- ◆ positive s_i is rightward, positive s_o is leftward
- ◆ positive y is upward

From Gauss's ray construction to the Gaussian lens formula



$$\frac{|y_i|}{y_o} = \frac{s_i}{s_o}$$

From Gauss's ray construction to the Gaussian lens formula

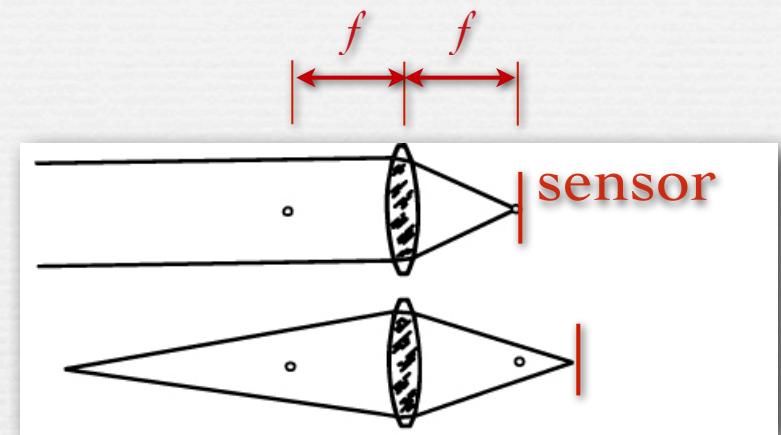


$$\frac{|y_i|}{y_o} = \frac{s_i}{s_o} \quad \text{and} \quad \frac{|y_i|}{y_o} = \frac{s_i - f}{f} \quad \dots$$

$$\frac{1}{s_o} + \frac{1}{s_i} = \frac{1}{f}$$

Changing the focus distance

- ◆ to focus on objects at different distances, move sensor relative to lens



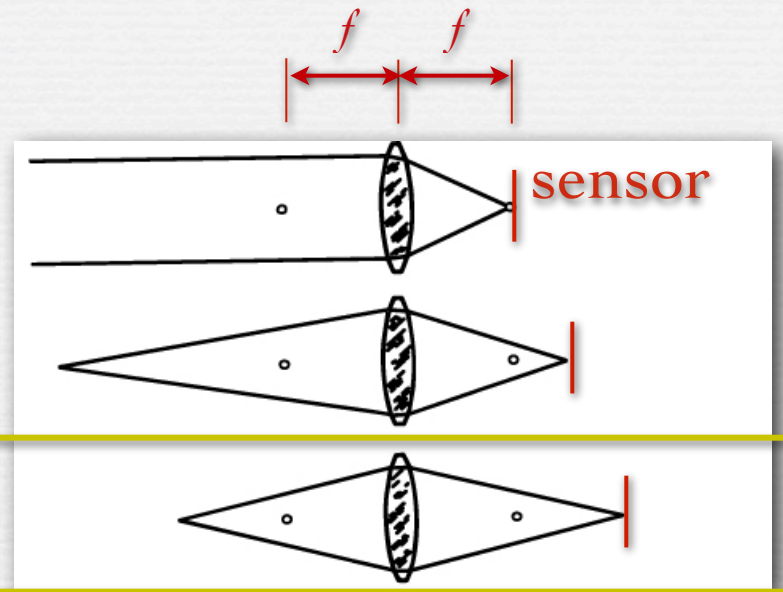
(FLASH DEMO)

<http://graphics.stanford.edu/courses/cs178/applets/gaussian.html>

$$\frac{1}{s_o} + \frac{1}{s_i} = \frac{1}{f}$$

Changing the focus distance

- ◆ to focus on objects at different distances, move sensor relative to lens



- ◆ at $s_o = s_i = 2f$ we have 1:1 imaging, because

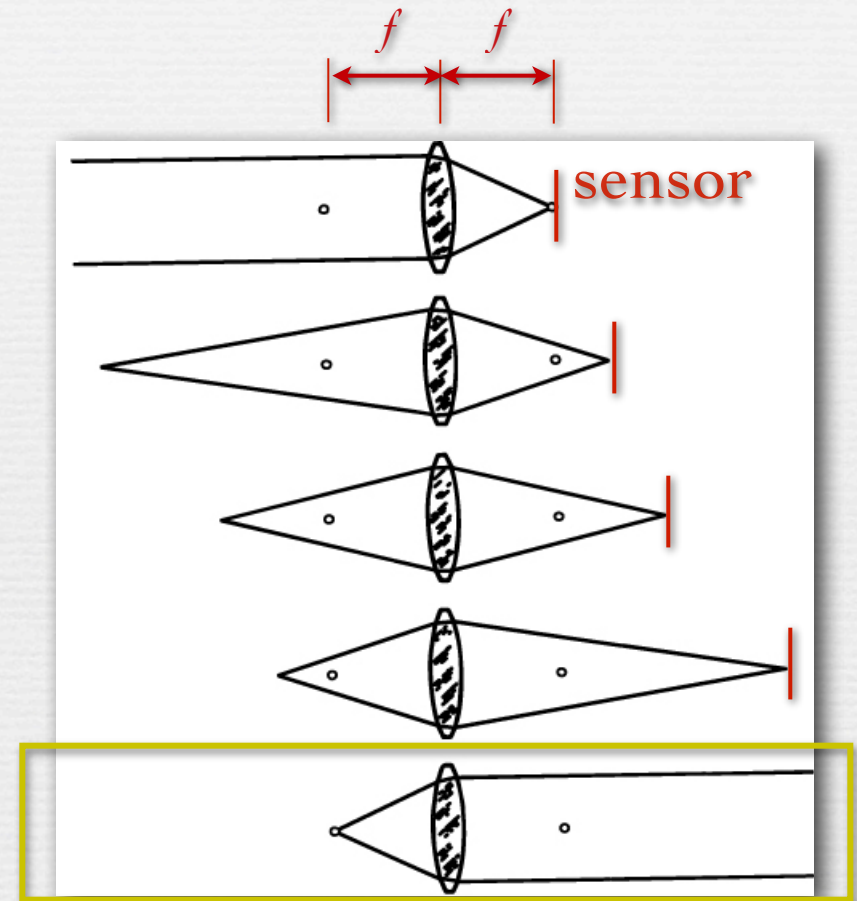
$$\frac{1}{2f} + \frac{1}{2f} = \frac{1}{f}$$

In 1:1 imaging, if the sensor is 36mm wide, an object 36mm wide will fill the frame.

$$\frac{1}{s_o} + \frac{1}{s_i} = \frac{1}{f}$$

Changing the focus distance

- ◆ to focus on objects at different distances, move sensor relative to lens
- ◆ at $s_o = s_i = 2f$ we have 1:1 imaging, because
$$\frac{1}{2f} + \frac{1}{2f} = \frac{1}{f}$$
- ◆ can't focus on objects closer to lens than its focal length f



$$\frac{1}{s_o} + \frac{1}{s_i} = \frac{1}{f}$$

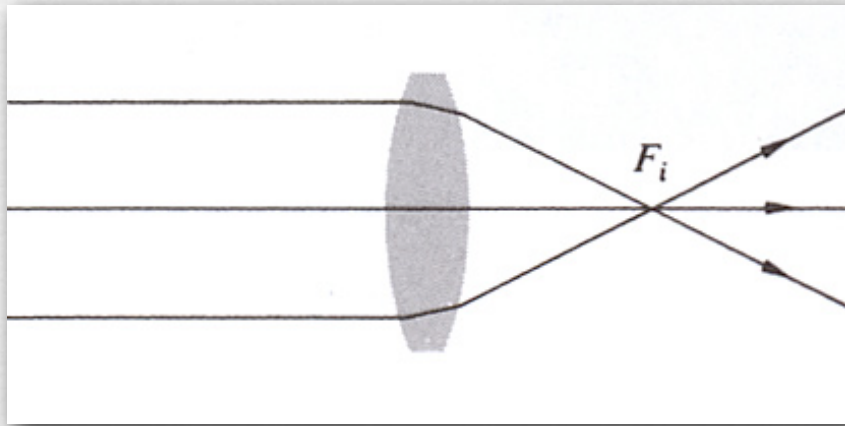
Recap

- ◆ approximations we sometimes make when analyzing lenses
 - *geometrical* optics instead of physical optics
 - *spherical* lenses instead of hyperbolic lenses
 - *thin lens* representation of thick optical systems
 - *paraxial* approximation of ray angles
- ◆ the Gaussian lens formula relates focal length f , object distance s_o , and image distance s_i
 - these settings, and sensor size, determine field of view
 - 1:1 imaging means $s_o = s_i$ and both are $2\times$ focal length
 - $s_o = f$ is the minimum possible object distance for a lens

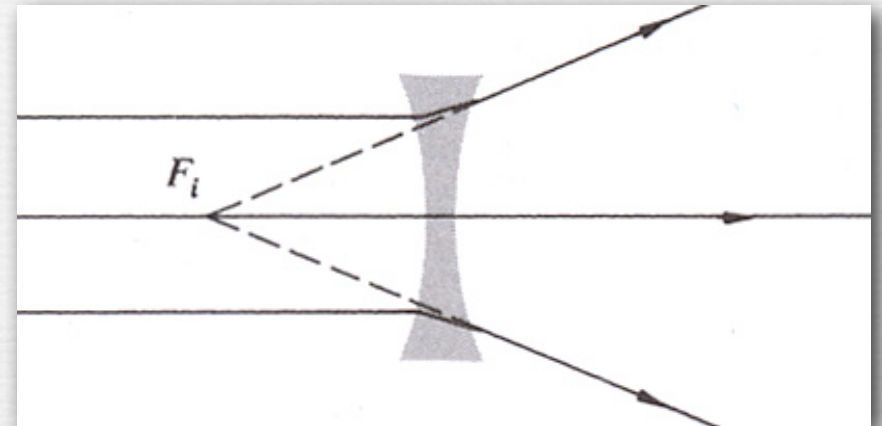
Questions?

Convex versus concave lenses

(Hecht)



rays from a convex lens converge

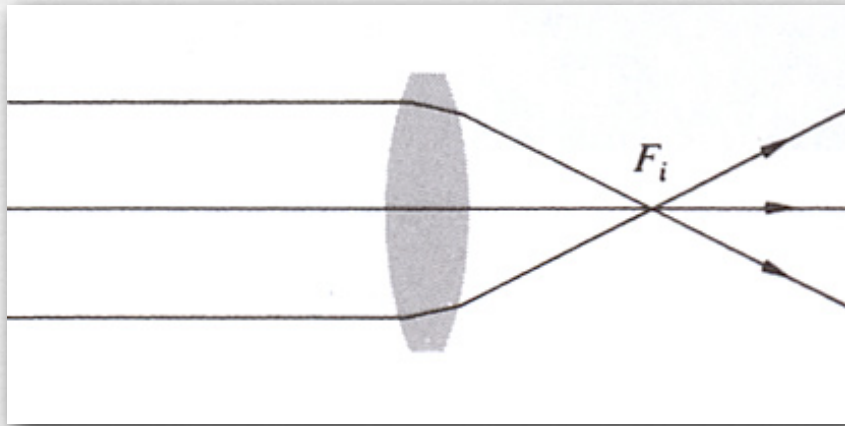


rays from a concave lens diverge

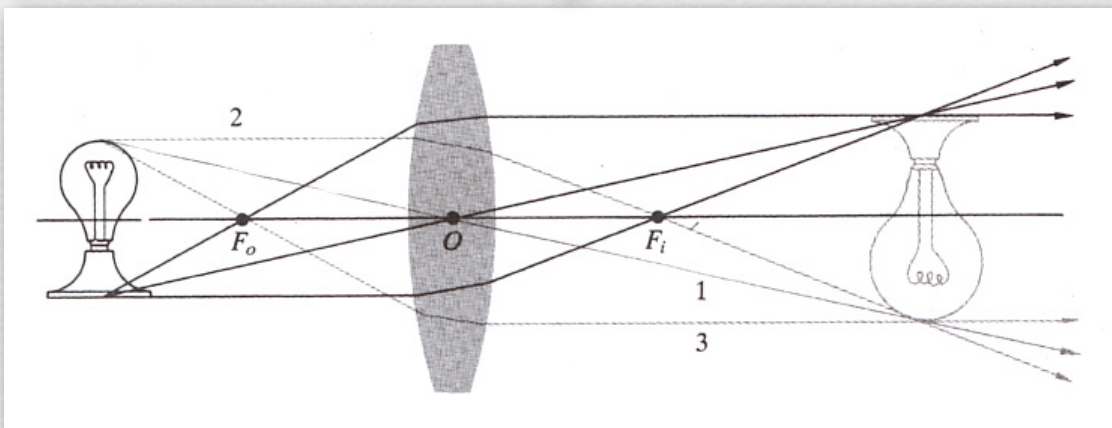
- ♦ positive focal length f means parallel rays from the left converge to a point on the right
- ♦ negative focal length f means parallel rays from the left converge to a point on the left (dashed lines above)

Convex versus concave lenses

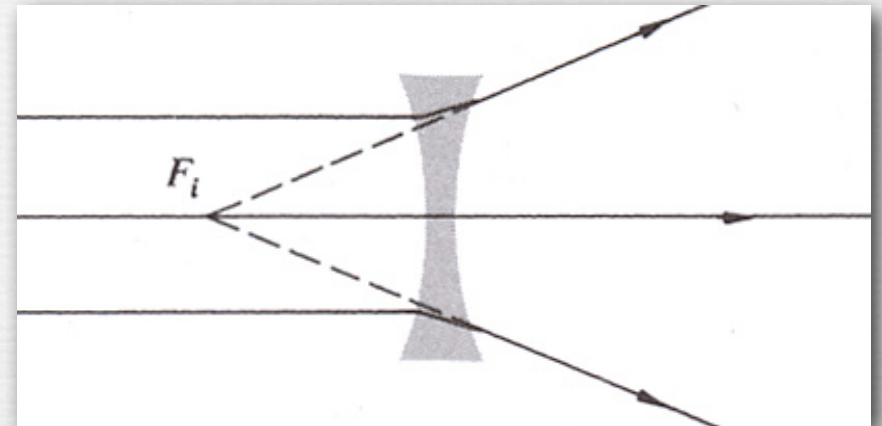
(Hecht)



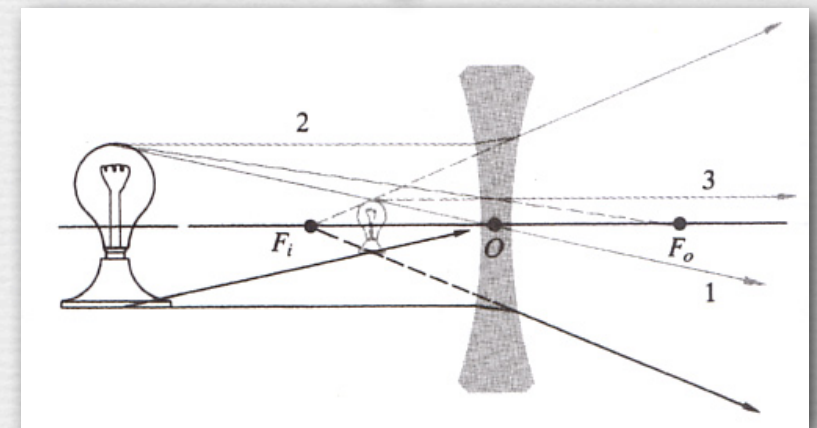
rays from a convex lens converge



...producing a real image

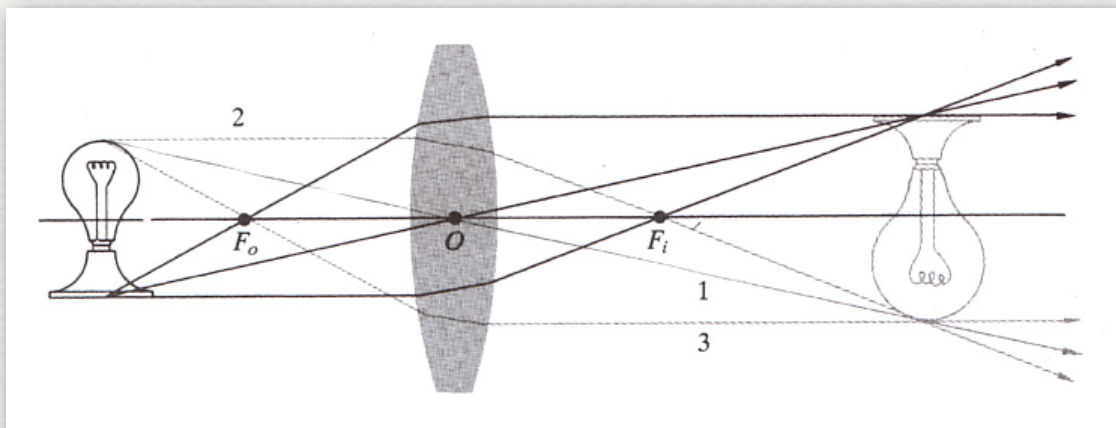
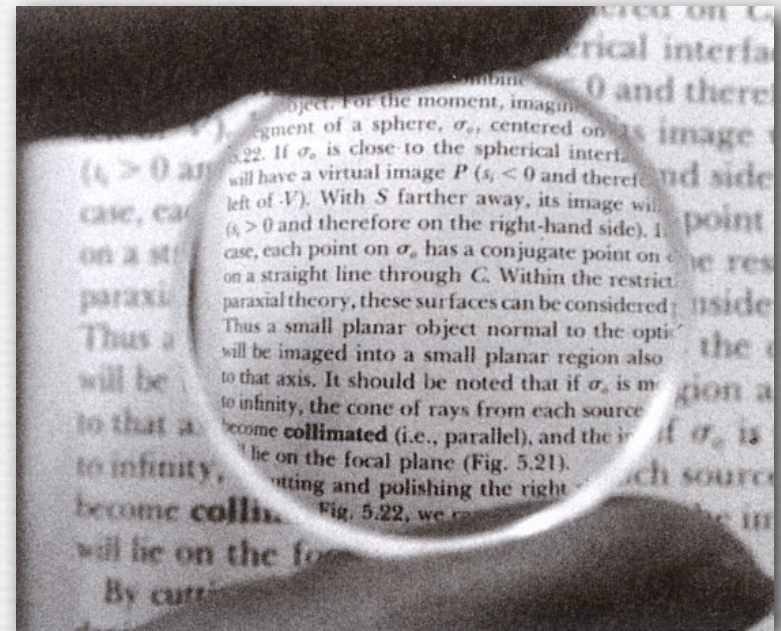
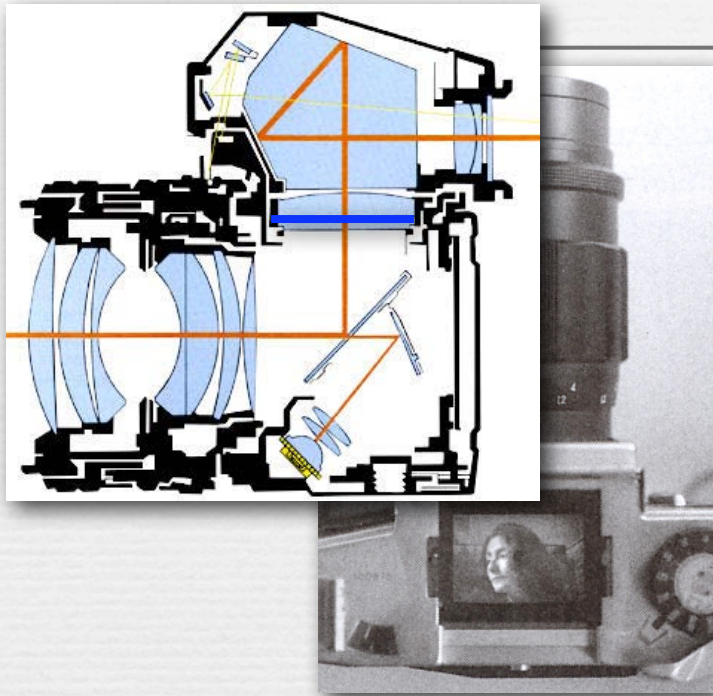


rays from a concave lens diverge

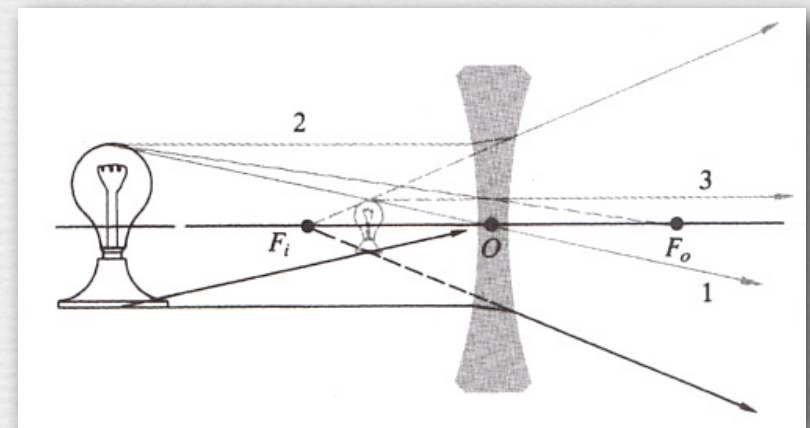


...producing a virtual image

Convex versus concave lenses



...producing a real image



...producing a virtual image

The power of a lens

$$P = \frac{1}{f}$$

units are meters⁻¹

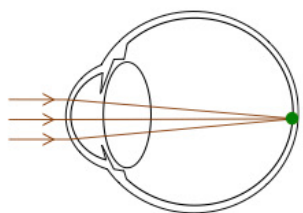
a.k.a. diopters

- ◆ my eyeglasses have the prescription
 - right eye: -0.75 diopters
 - left eye: -1.00 diopters

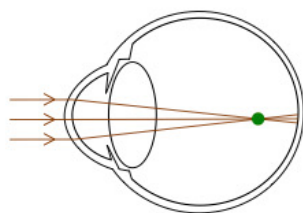
Q. What's wrong with me?

A. Myopia (nearsightedness)

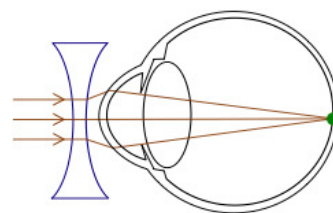
(Pamplona)



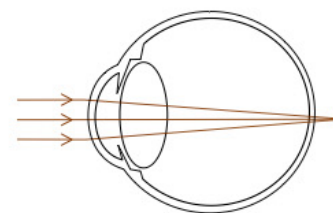
(a) Perfect eye



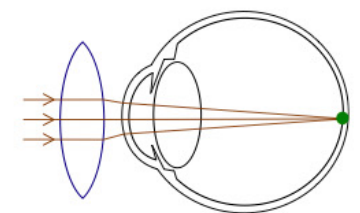
(b) Myopia



(c) Corrected Myopia



(d) Hyperopia



(e) Corrected Hyperopia

Combining two lenses

- ◆ using focal lengths

$$\frac{1}{f_{tot}} = \frac{1}{f_1} + \frac{1}{f_2}$$

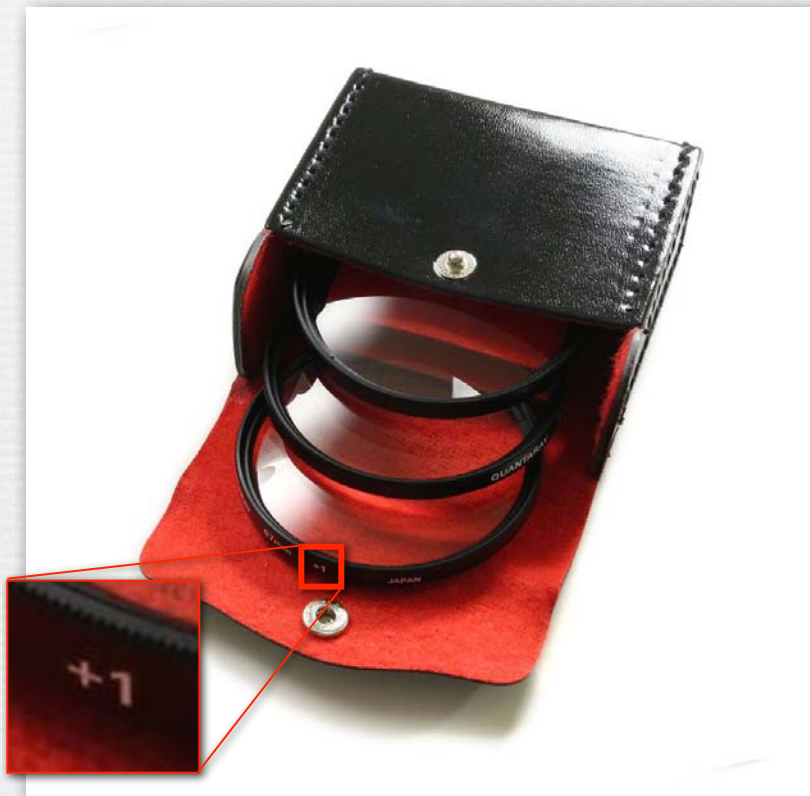
- ◆ using diopters

$$P_{tot} = P_1 + P_2$$

- ◆ example

$$\frac{1}{200mm} + \frac{1}{500mm} = \frac{1}{143mm} \quad \text{-or-} \quad 5.0 + 2.0 = 7.0 \text{ diopters}$$

Close-up filters



(wikipedia)

- ◆ screw on to end of lens
- ◆ power is designated in diopters (usually)

Close-up filters

Panasonic 45-200



+



- ◆ changes longest focal length from 200mm to 143mm

$$\frac{1}{200\text{mm}} + \frac{1}{500\text{mm}} = \frac{1}{143\text{mm}} \quad \text{-or-} \quad 5.0 + 2.0 = 7.0 \text{ diopters}$$

Close-up filters

- ◆ for a fixed image distance, it reduces the object distance
 - at $f=200\text{mm}$, this len's minimum object distance $s_o = 1000\text{mm}$
 - at these settings, its effective image distance must be

$$s_i = \frac{1}{\frac{1}{f} - \frac{1}{s_o}} = \frac{1}{\frac{1}{200\text{mm}} - \frac{1}{1000\text{mm}}} = 250\text{mm}$$

- with the closeup filter and the same settings of focal length and image distance, the in-focus object distance becomes

$$s_o = \frac{1}{\frac{1}{f} - \frac{1}{s_i}} = \frac{1}{\frac{1}{143\text{mm}} - \frac{1}{250\text{mm}}} = 334\text{mm}$$

3x
closer!

Close-up filters



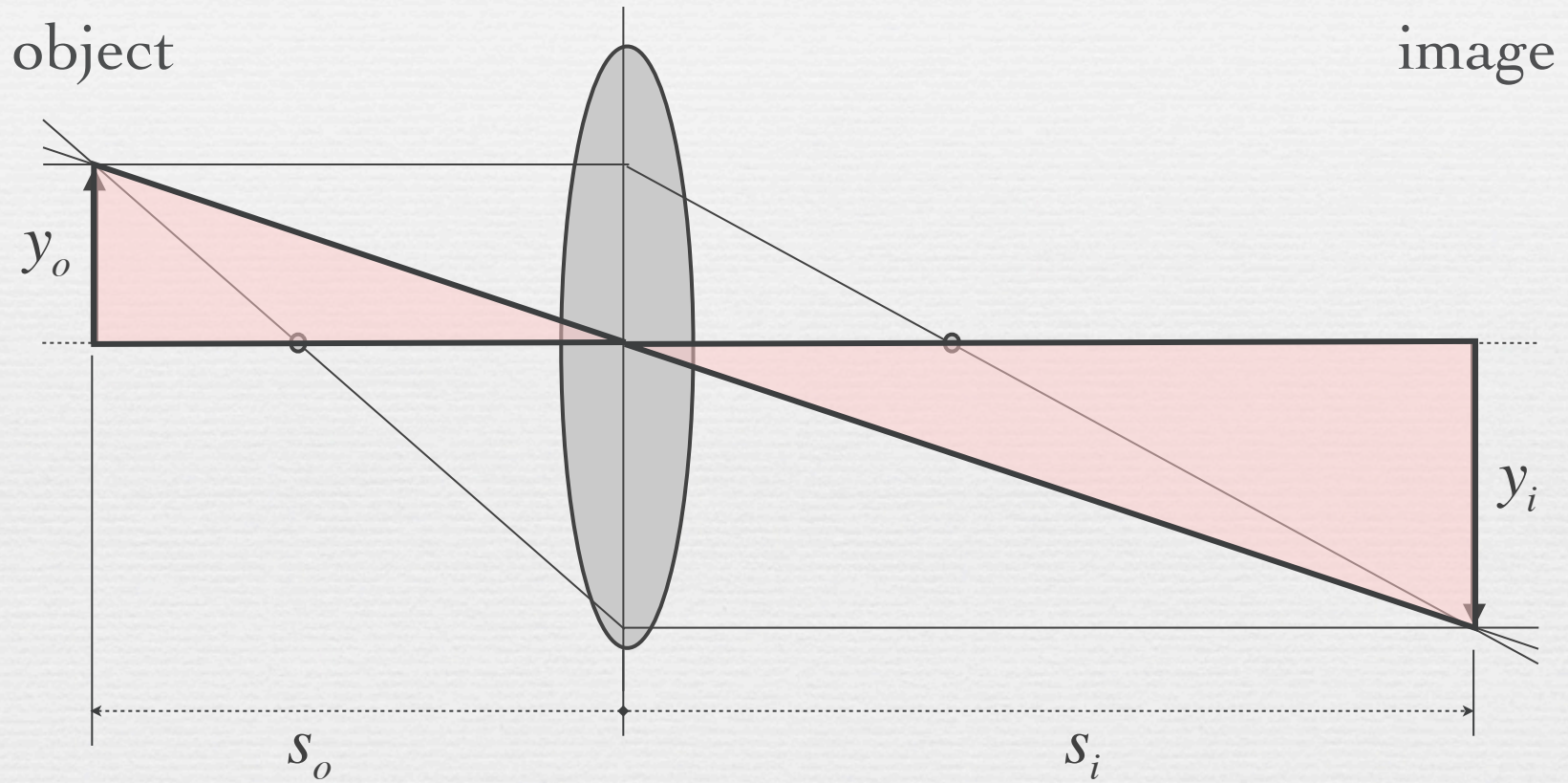
200mm lens &
no closeup filter
 $s_o = 1000\text{mm}$



200mm lens &
500D closeup filter
 $s_o = 334\text{mm}$

poor man's
macro lens

Magnification



$$M_T \triangleq \frac{y_i}{y_o} = -\frac{s_i}{s_o}$$

Close-up filters



200mm lens &
no closeup filter
 $s_o = 1000\text{mm}$

$$M_T = -\frac{s_i}{s_o} = \frac{250}{1000} = -1:4$$

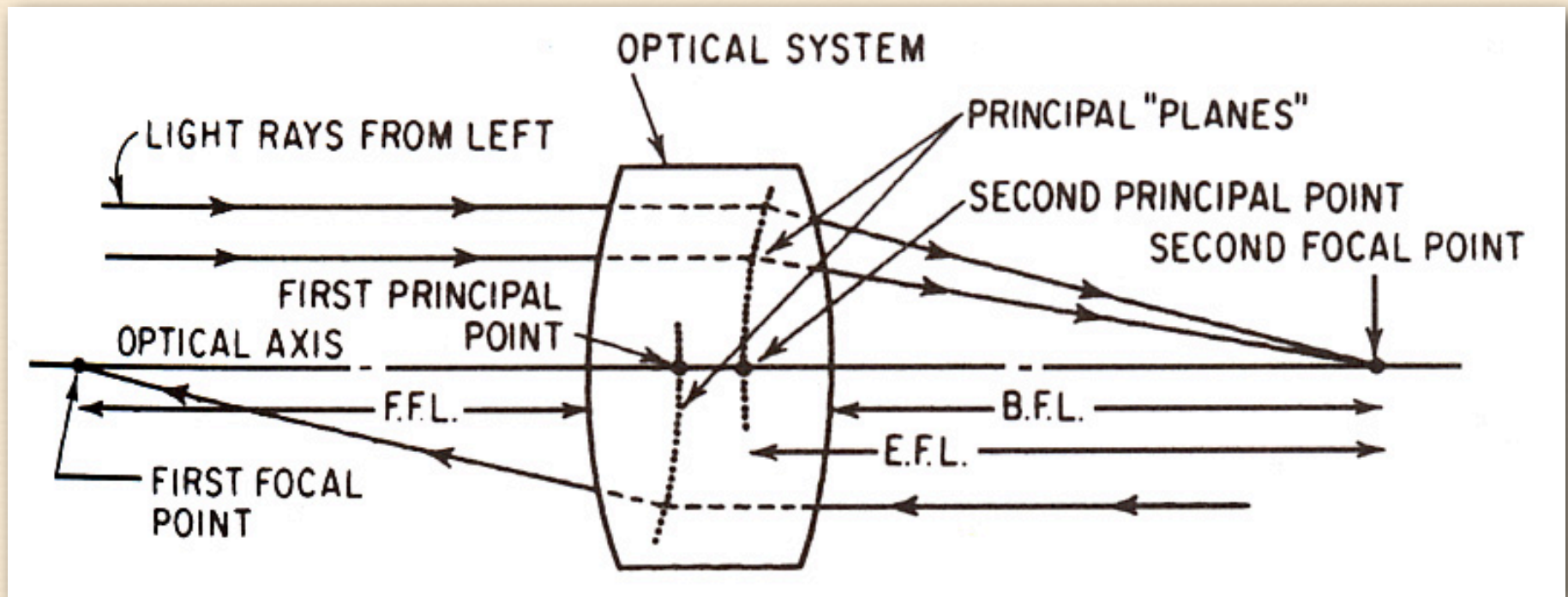


200mm lens &
500D closeup filter
 $s_o = 334\text{mm}$

$$M_T = -\frac{s_i}{s_o} = \frac{250}{334} = -3:4$$

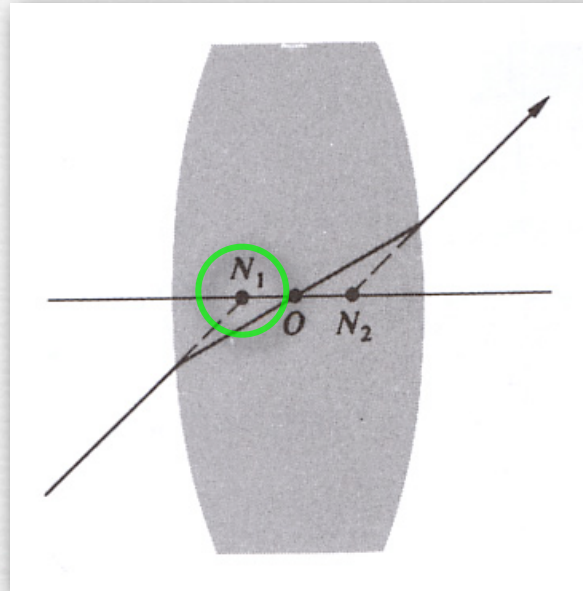
Thick lenses

- ♦ an optical system may contain many lenses, but can be characterized by a few numbers



(Smith)

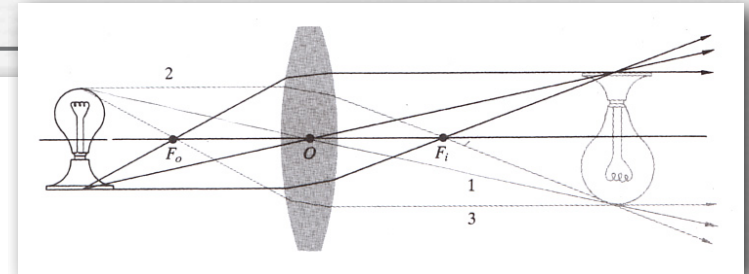
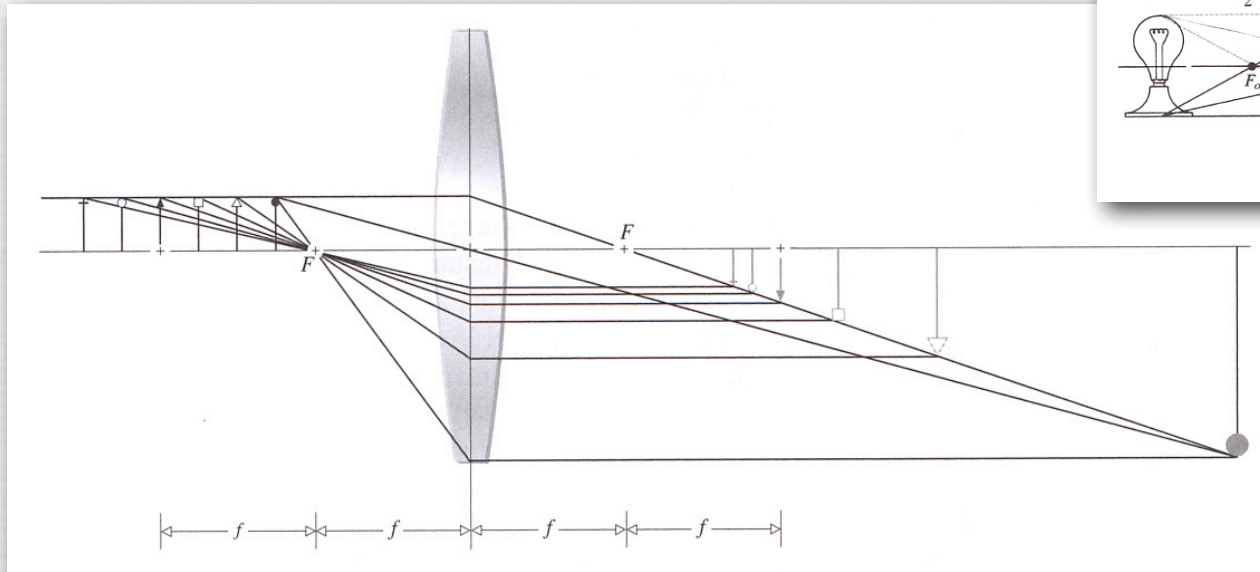
Center of perspective



(Hecht)

- in a thin lens, the *chief ray* from a point traverses the lens (through its optical center) without changing direction
- in a thick lens, the intersections of this ray with the optical axis are called the *nodal points*
- for a lens in air, these coincide with the *principal points*
- the first nodal point is the *center of perspective*

Lenses perform a 3D perspective transform



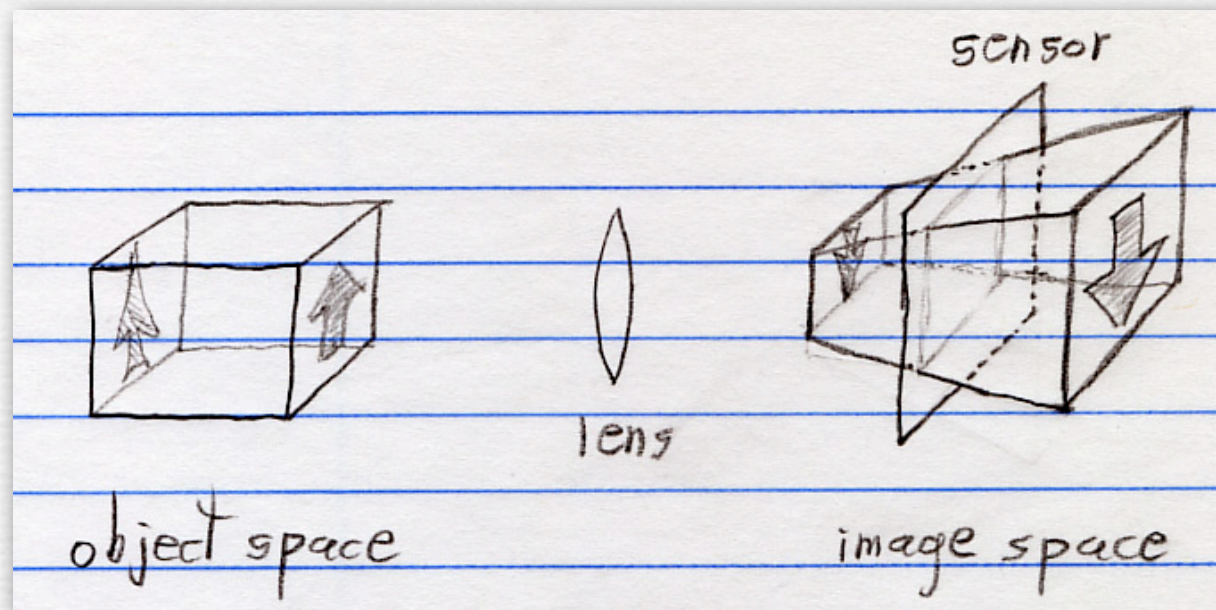
(FLASH DEMO)

<http://graphics.stanford.edu/courses/cs178/applets/thinlens.html>

(Hecht)

- ◆ lenses transform a 3D object to a 3D image; the sensor extracts a 2D slice from that image
- ◆ as an object moves linearly (in Z), its image moves non-proportionately (in Z)
- ◆ as you move a lens linearly relative to the sensor, the in-focus object plane moves non-proportionately
- ◆ as you refocus a camera, the image changes size !

Lenses perform a 3D perspective transform (contents of whiteboard)



- ◆ a cube in object space is transformed by a lens into a 3D frustum in image space, with the orientations shown by the arrows
- ◆ in computer graphics this transformation is modeled as a 4×4 matrix multiplication of 3D points expressed in 4D homogenous coordinates
- ◆ in photography a sensor extracts a 2D slice from the 3D frustum; on this slice some objects will be sharply focused; others may be blurry

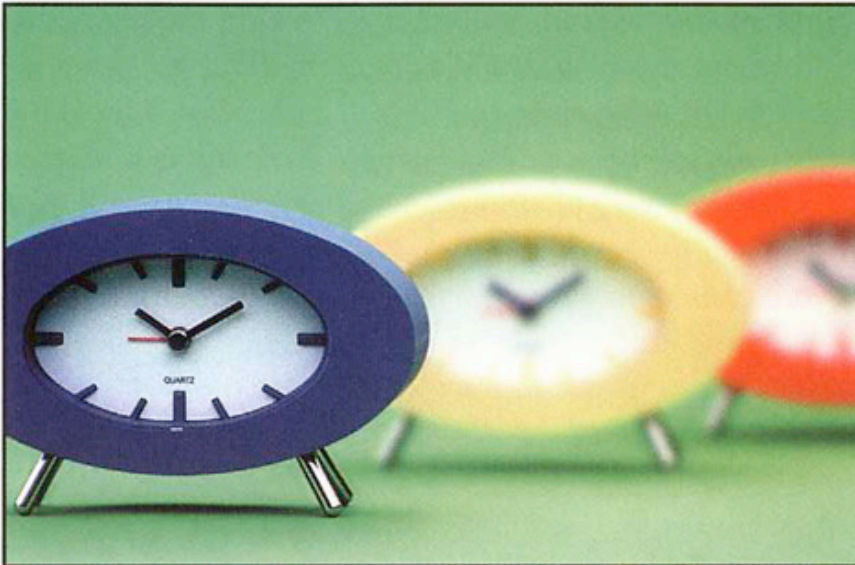
Recap

- ◆ more implications of the Gaussian lens formula
 - convex lenses make real images; concave make virtual images
 - the power of a lens (in diopters) is 1 over its focal length
 - when combining two lenses, add their powers
 - adding a closeup filter allows a smaller object distance
 - changing object and image distances changes magnification
- ◆ lenses perform a 3D perspective transform of object space
 - an object's apparent size is inversely proportional to its distance
 - linear lens motions move the in-focus plane non-linearly
 - focusing a lens changes the image size (slightly)

Questions?

Depth of field

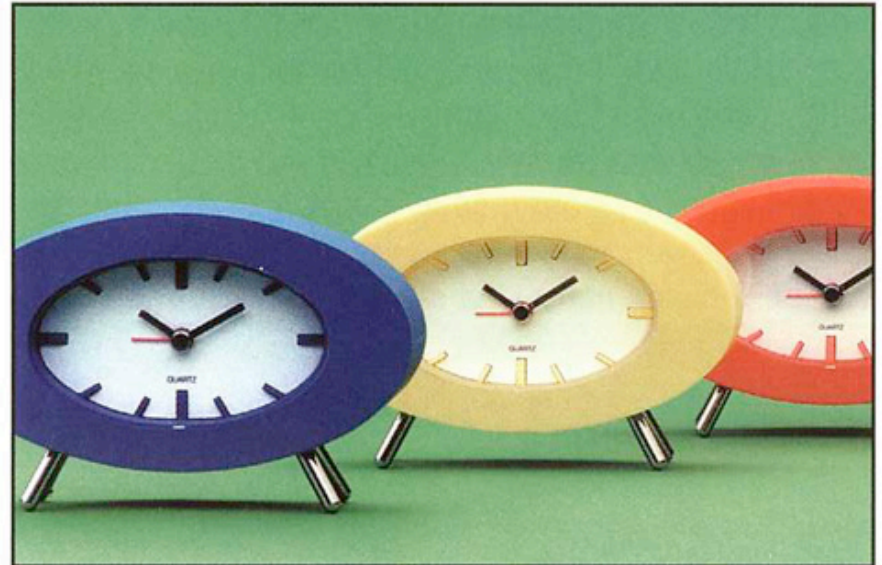
LESS DEPTH OF FIELD



Wider aperture



MORE DEPTH OF FIELD



Smaller aperture

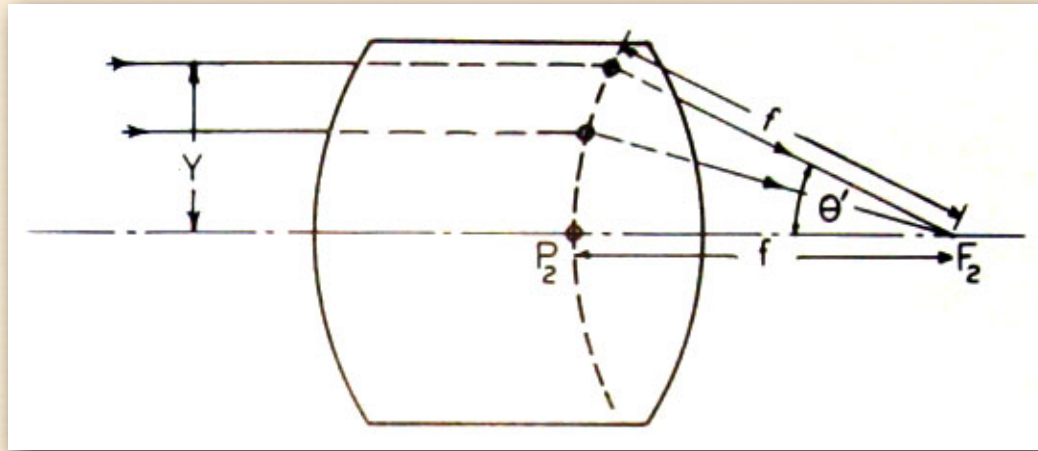


(London)

$$N = \frac{f}{A}$$

- ◆ lower N means a wider aperture and less depth of field

How low can N be?



(Kingslake)

- ◆ principal planes are the paraxial approximation of a spherical “equivalent refracting surface”

$$N = \frac{1}{2 \sin \theta'}$$

- ◆ lowest possible N in air is f/0.5
- ◆ lowest N I've seen in an SLR is f/1.0



Canon EOS 50mm f/1.0
(discontinued)

Cinematography by candlelight



Stanley Kubrick,
Barry Lyndon,
1975

- ◆ Zeiss 50mm f/0.7 Planar lens
 - originally developed for NASA's Apollo missions
 - very shallow depth of field in closeups (small object distance)

Cinematography by candlelight

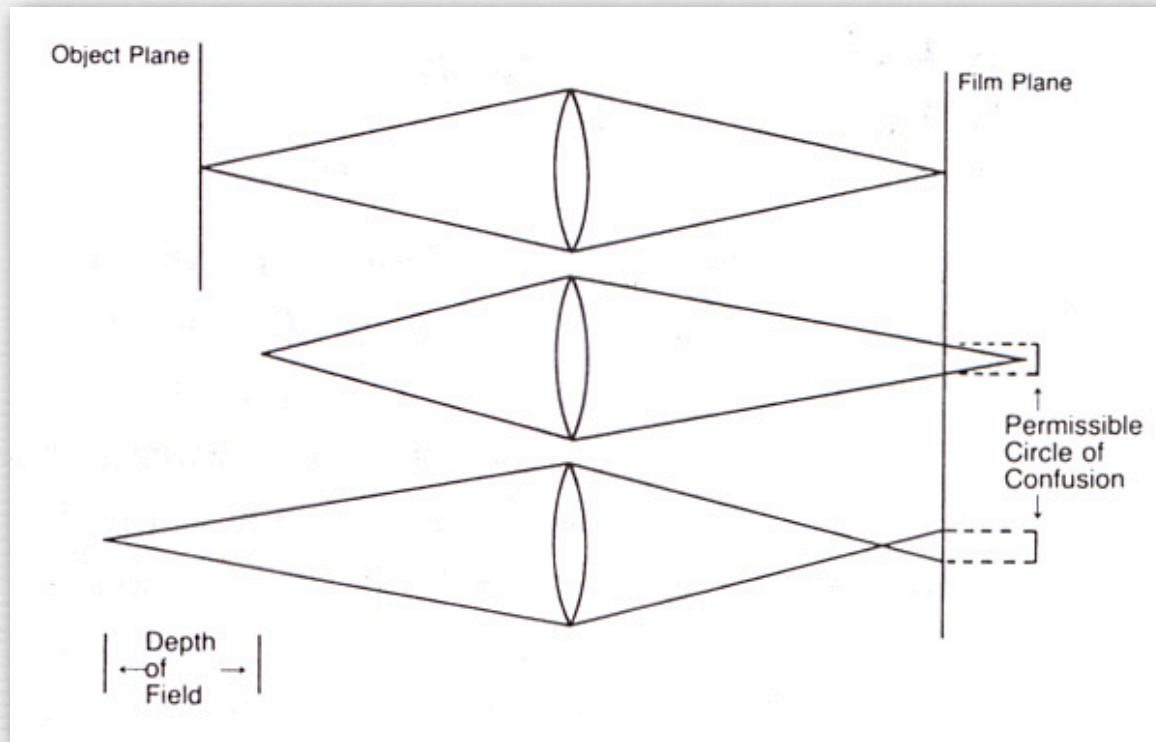


Stanley Kubrick,
Barry Lyndon,
1975



- ◆ Zeiss 50mm f/0.7 Planar lens
 - originally developed for NASA's Apollo missions
 - very shallow depth of field in closeups (small object distance)

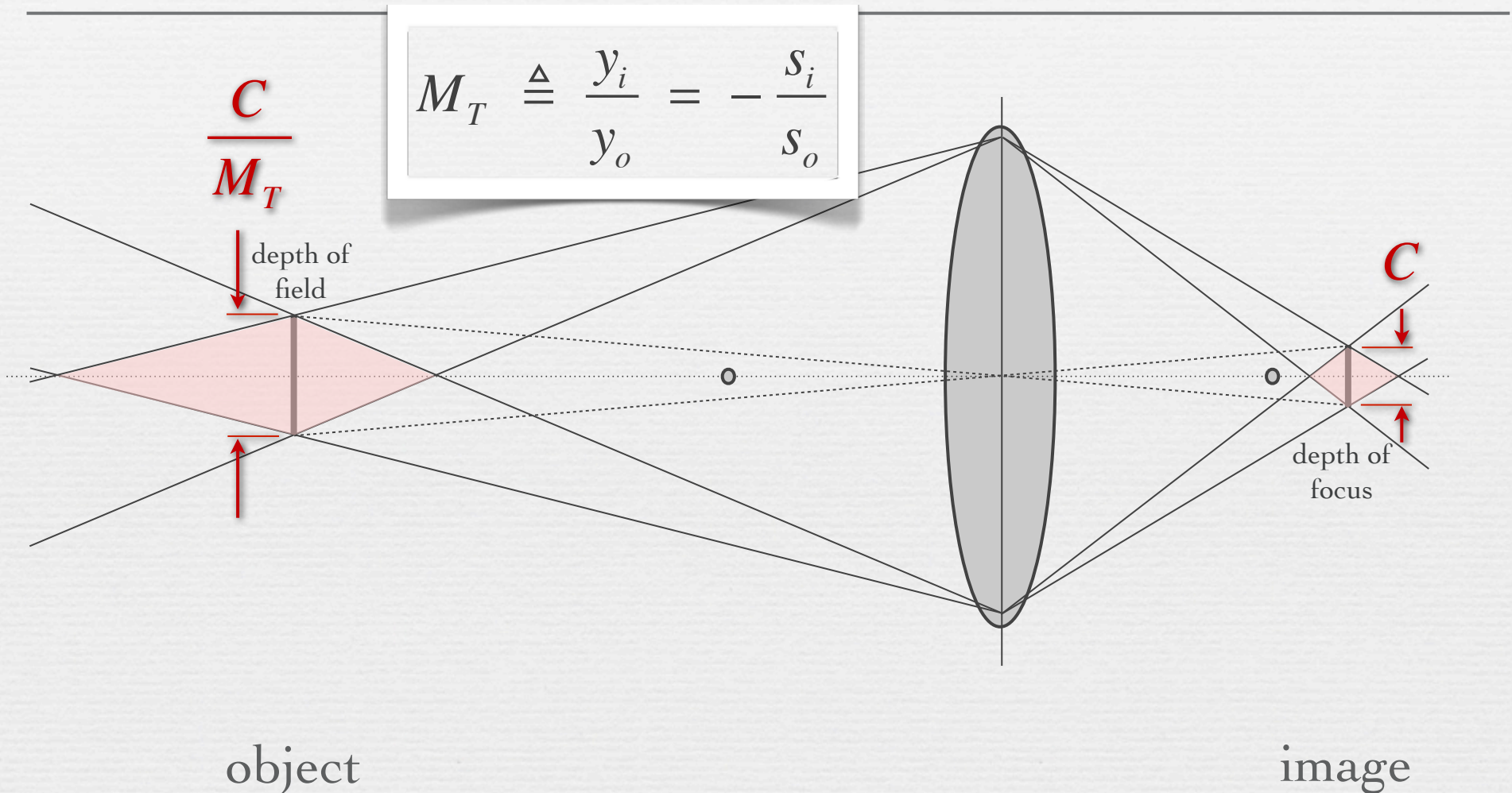
Circle of confusion (C)



to derive C for a given situation, start from the smallest visual angle we can detect; we'll cover this later in the course

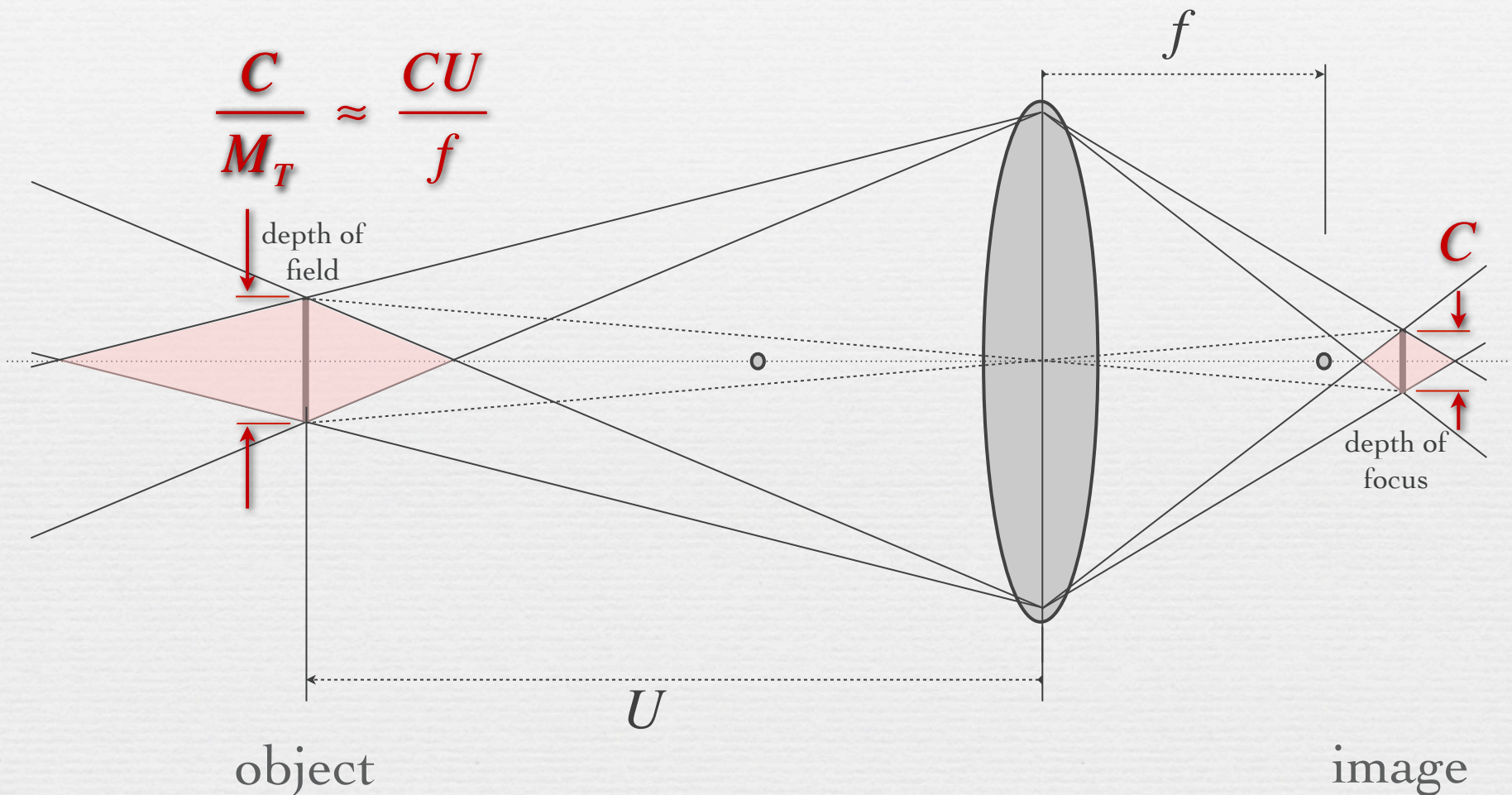
- ◆ C depends on sensing medium, reproduction medium, viewing distance, human vision, ...
 - for print from 35mm film, 0.02mm (on negative) is typical
 - for high-end SLR, 6 μ is typical (1 pixel)
 - larger if downsizing for web, or lens is poor

Depth of field formula



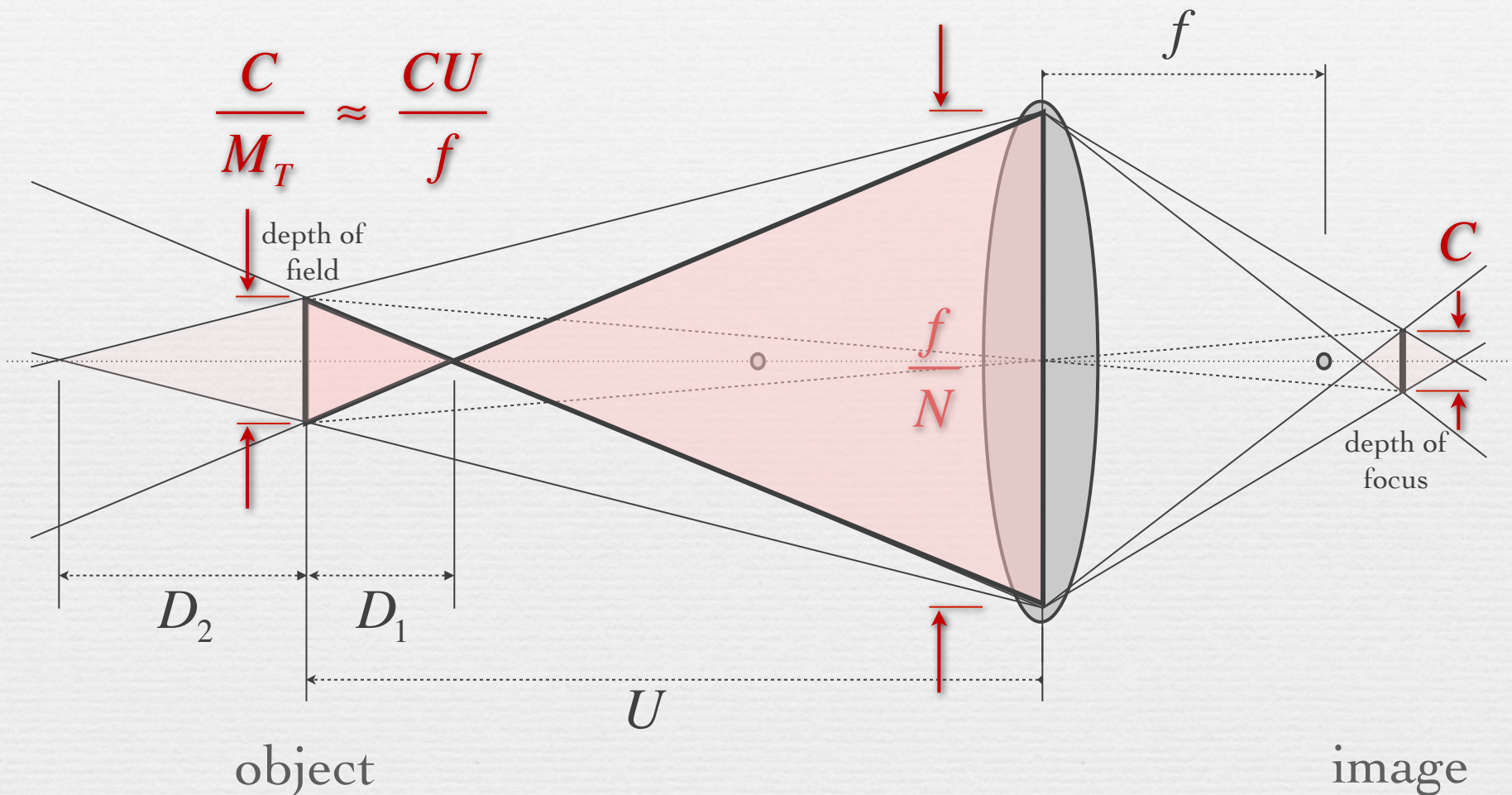
- ◆ DoF is asymmetrical around the in-focus object plane
- ◆ conjugate in object space is typically bigger than C

Depth of field formula



- ◆ DoF is asymmetrical around the in-focus object plane
- ◆ conjugate in object space is typically bigger than C

Depth of field formula



$$\frac{D_1}{CU / f} = \frac{U - D_1}{f / N} \dots D_1 = \frac{NCU^2}{f^2 + NCU} \quad D_2 = \frac{NCU^2}{f^2 - NCU}$$

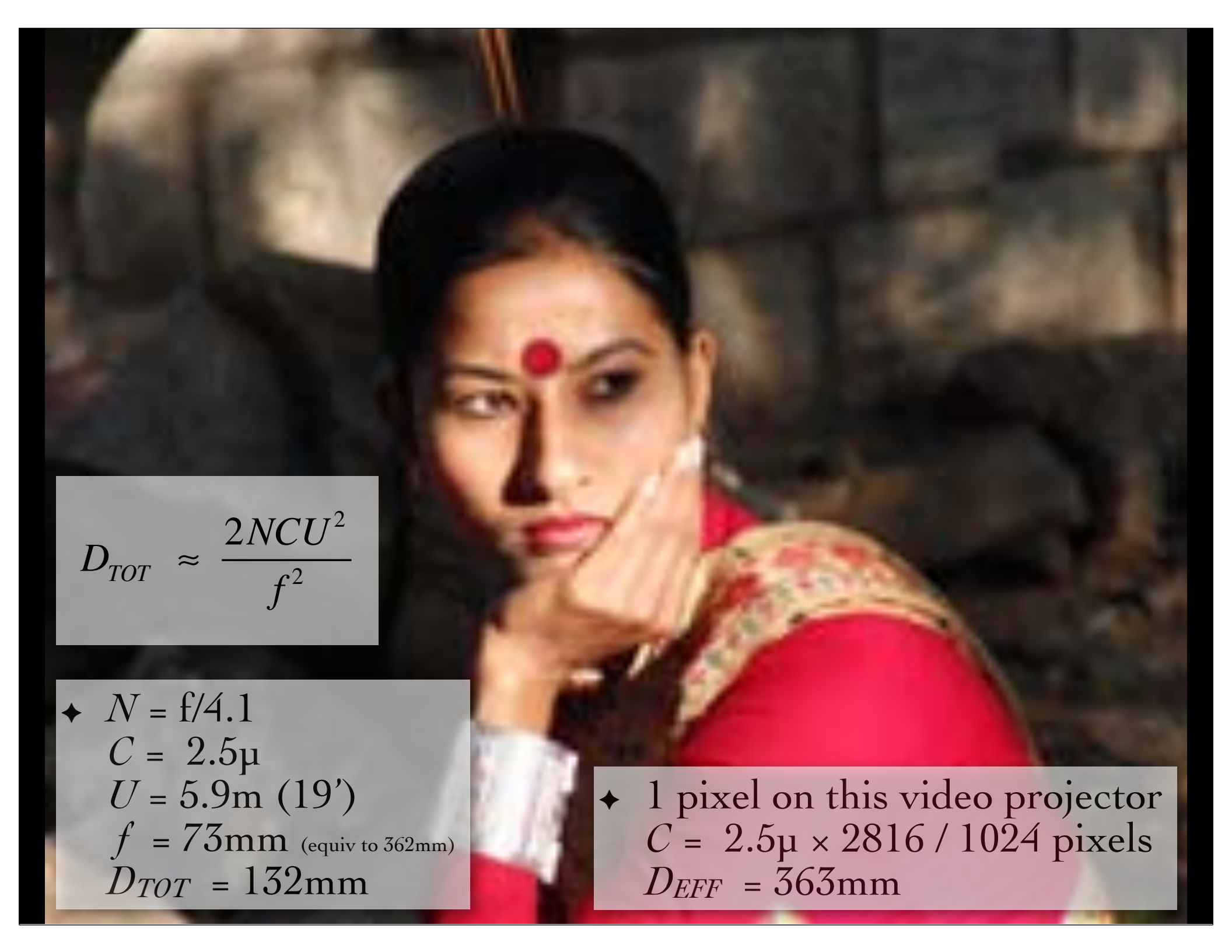
Depth of field formula

$$D_{TOT} = D_1 + D_2 = \frac{2NCU^2 f^2}{f^4 - N^2 C^2 U^2}$$

- ◆ $N^2 C^2 U^2$ can be ignored when conjugate of circle of confusion is small relative to the aperture

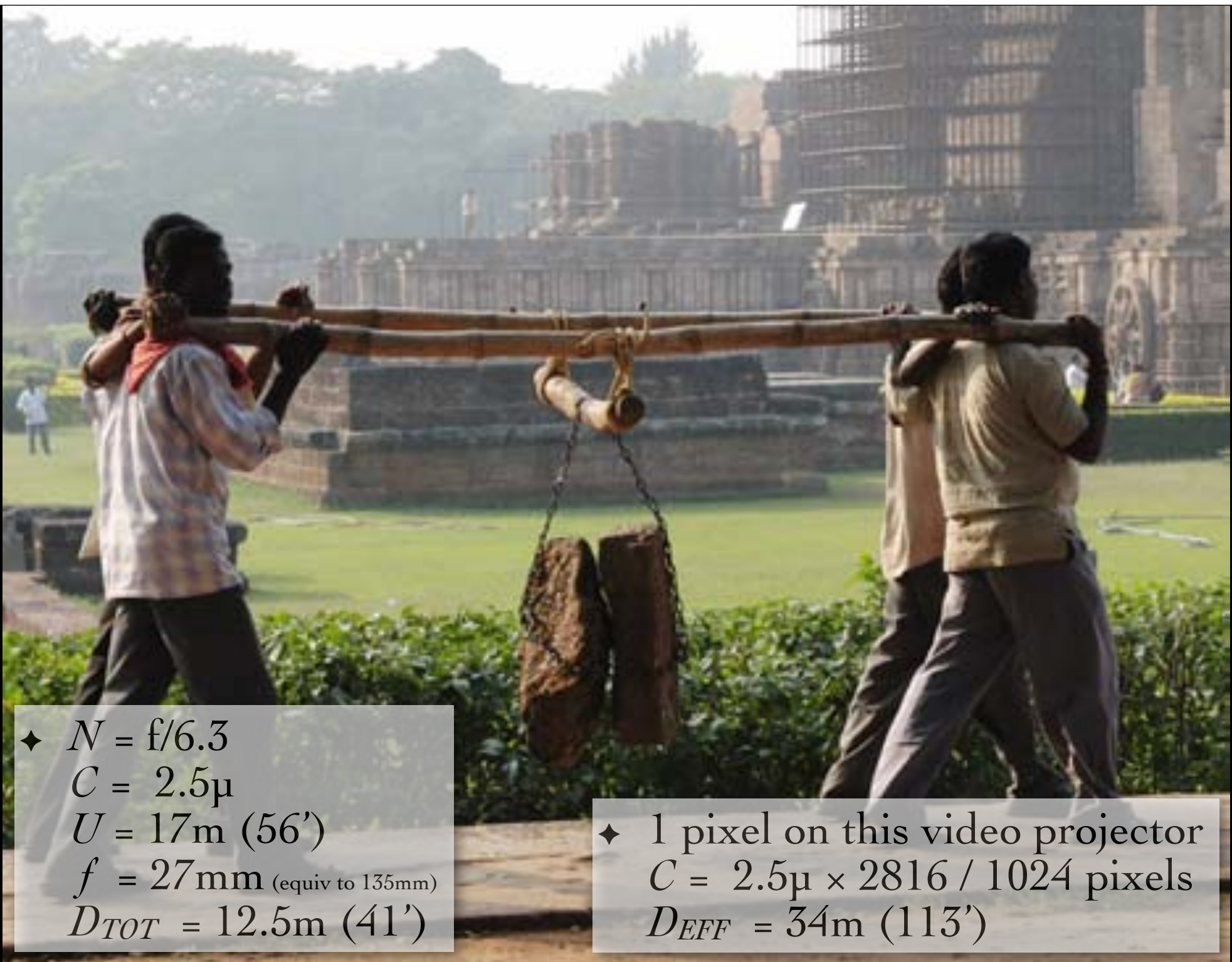
$$D_{TOT} \approx \frac{2NCU^2}{f^2}$$

- ◆ where
 - N is F-number of lens
 - C is circle of confusion (on image)
 - U is distance to in-focus plane (in object space)
 - f is focal length of lens


$$D_{TOT} \approx \frac{2NCU^2}{f^2}$$


- ◆ $N = f/4.1$
- $C = 2.5\mu$
- $U = 5.9\text{m (19')}$
- $f = 73\text{mm (equiv to 362mm)}$
- $D_{TOT} = 132\text{mm}$

- ◆ 1 pixel on this video projector
 $C = 2.5\mu \times 2816 / 1024$ pixels
 $D_{EFF} = 363\text{mm}$



◆ $N = f/6.3$
 $C = 2.5\mu$
 $U = 17\text{m (56')}$
 $f = 27\text{mm (equiv to 135mm)}$
 $D_{TOT} = 12.5\text{m (41')}$

◆ 1 pixel on this video projector
 $C = 2.5\mu \times 2816 / 1024$ pixels
 $D_{EFF} = 34\text{m (113')}$



◆ $N = f/5.6$
 $C = 6.4\mu$
 $U = 0.7\text{m}$
 $f = 105\text{mm}$
 $D_{TOT} = 3.2\text{mm}$

◆ 1 pixel on this video projector
 $C = 6.4\mu \times 5616 / 1024$ pixels
 $D_{EFF} = 17.5\text{mm}$



Canon MP-E
65mm 5:1 macro



M.S.

◆ $N = f/2.8$
 $C = 6.4\mu$
 $U = 78\text{mm}$
 $f = 65\text{mm}$

(use $N' = (1+M_T)N$ at short conjugates ($M_T=5$ here)) = $f/16$

$D_{TOT} = 0.29\text{mm}!$

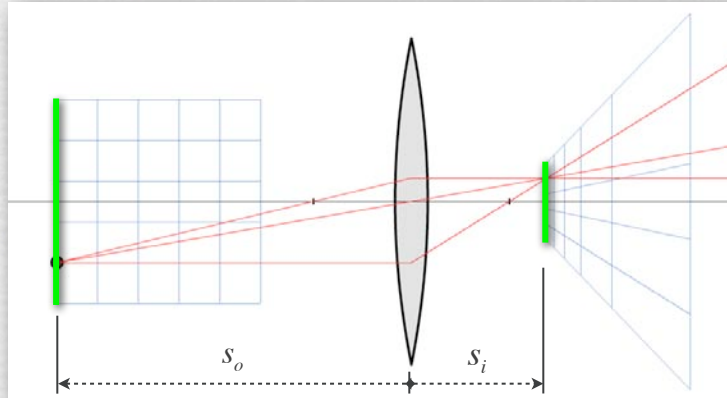
(Mikhail Shlemov)

Sidelight: macro lenses

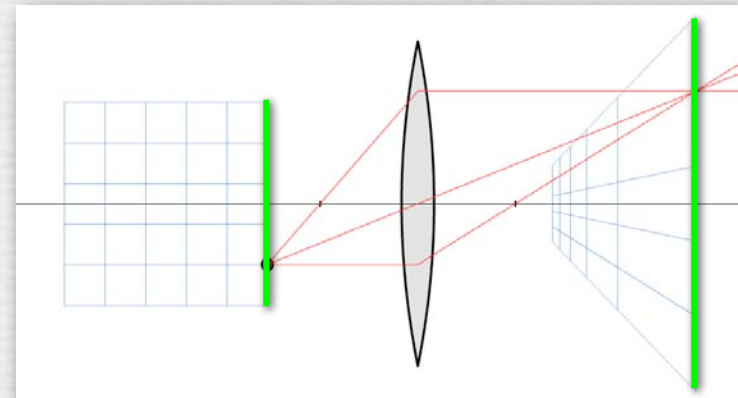
$$\frac{1}{s_o} + \frac{1}{s_i} = \frac{1}{f}$$



Q. How can the Casio EX-F1 at 73mm and the Canon MP-E 65mm macro, which have similar f 's, have such different focusing distances?



normal



macro

- ◆ A. Because macro lenses are built to allow long s_i
 - this changes s_o , which changes magnification $M_T \triangleq -s_i / s_o$
 - macro lenses are also well corrected for aberrations at short s_o

Extension tube: fits between camera and lens, converts a normal lens to a macro lens



- ◆ toilet paper tube, black construction paper, masking tape
- ◆ camera hack by Katie Dektar (CS 178, 2009)

Extension tubes versus close-up filters



Canon 25mm

In past years students have asked about "close-focus" settings on some standard zoom lenses. These indeed allow macro shots to be captured, but such lenses are not well corrected for operating at close object distances, so they will produce mediocre-quality macro images.



Canon $f = 500\text{mm}$

- ◆ both allow closer focusing, hence greater magnification
- ◆ both degrade image quality relative to a macro lens
- ◆ extension tubes work best with wide-angle lenses; close-up filters work best with telephoto lenses
- ◆ extension tubes raise F-number, reducing light
- ◆ need different close-up filter for each lens filter diameter

Extension tubes versus close-up filters versus teleconverters



Canon 25mm



Canon $f = 500\text{mm}$



Nikon 1.4x

- ◆ a teleconverter fits between the camera and lens, like an extension tube
- ◆ they increase f , narrowing FOV & increasing magnification, but they don't change the focusing range
- ◆ like extension tubes, they raise F-number, reducing light, and they are awkward to add or remove
- ◆ see <http://www.cambridgeincolour.com/tutorials/macro-extension-tubes-closeup.htm>

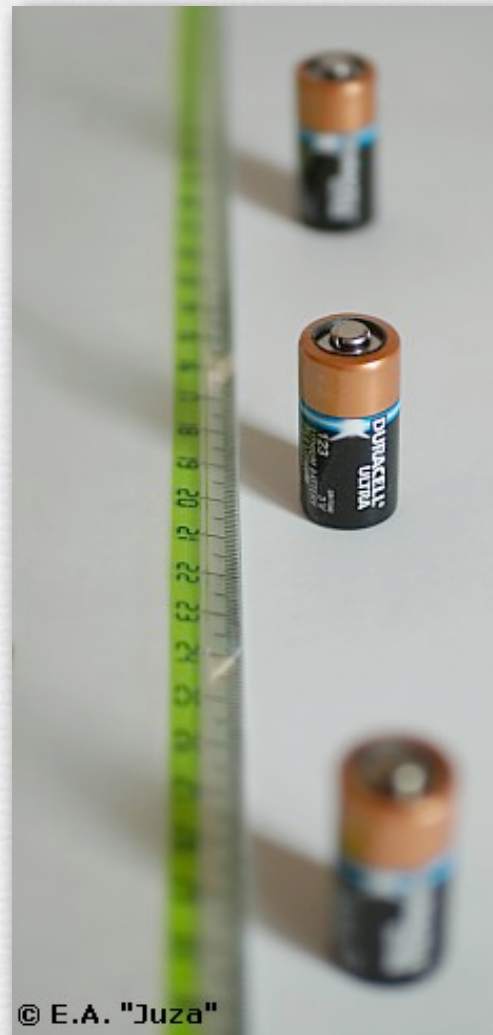
DoF is linear with F-number

$$D_{TOT} \approx \frac{2NCU^2}{f^2}$$

(FLASH DEMO)

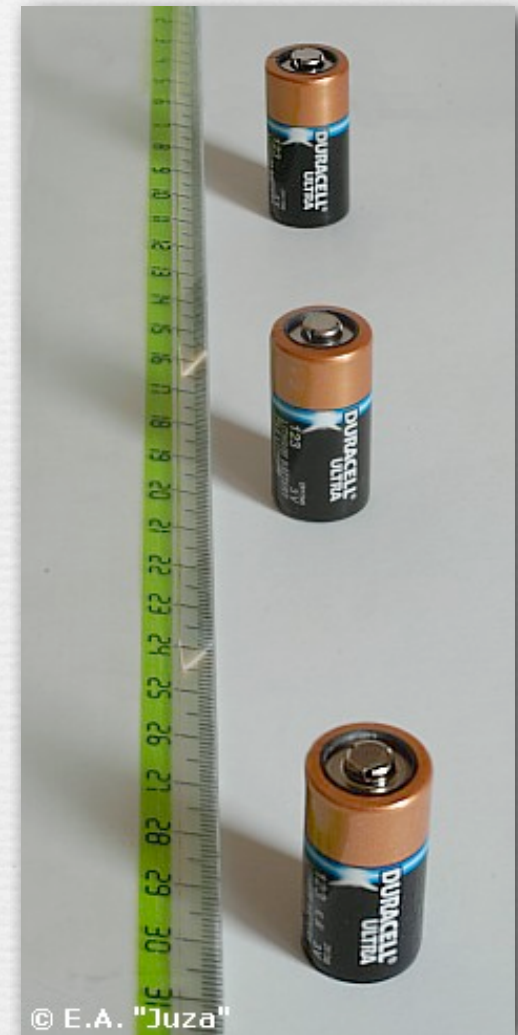
<http://graphics.stanford.edu/courses/cs178/applets/dof.html>

(juzaphoto.com)



© E.A. "Juza"

f/2.8



© E.A. "Juza"

f/32

DoF is quadratic with subject distance

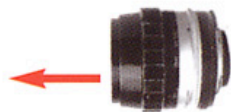
$$D_{TOT} \approx \frac{2NCU^2}{f^2}$$

(FLASH DEMO)

<http://graphics.stanford.edu/courses/cs178/applets/dof.html>



Closer to subject



3 feet



Farther from subject



10 feet

(London)

Hyperfocal distance



- ◆ the back depth of field

$$D_2 = \frac{NCU^2}{f^2 - NCU}$$

- ◆ becomes infinite if

$$U \geq \frac{f^2}{NC} \triangleq H$$

- ◆ In that case, the front depth of field becomes

$$D_1 = \frac{NCU^2}{f^2 + NCU} = \frac{H}{2}$$

- ◆ so if I had focused at 32m, everything from 16m to infinity would be in focus on a video projector, including the men at 17m

- ◆ $N = f/6.3$
 $C = 2.5\mu \times 2816 / 1920$ pixels
 $U = 17\text{m}$ (56')
 $f = 27\text{mm}$ (equiv to 135mm)
 $D_{TOT} = 18.3\text{m}$ on HD projector
 $H = 31.6\text{m}$ (104')

(FLASH DEMO)

<http://graphics.stanford.edu/courses/cs178/applets/dof.html>

DoF is inverse quadratic with focal length

$$D_{TOT} \approx \frac{2NCU^2}{f^2}$$

(FLASH DEMO)

<http://graphics.stanford.edu/courses/cs178/applets/dof.html>



Longer focal length



180mm



Shorter focal length



50mm

(London)

Q. Does sensor size affect DoF?

$$D_{TOT} \approx \frac{2NCU^2}{f^2}$$

- ◆ as sensor shrinks, lens focal length f typically shrinks to maintain a comparable field of view
- ◆ as sensor shrinks, pixel size C typically shrinks to maintain a comparable number of pixels in the image
- ◆ thus, depth of field D_{TOT} increases linearly with decreasing sensor size on consumer cameras
- ◆ this is why amateur cinematographers are drawn to SLRs
 - their chips are larger than even pro-level video camera chips
 - so they provide unprecedented control over depth of field



Vincent Laforet, Nocturne (2009)

Canon 1D Mark IV

DoF and the dolly-zoom

- ◆ if we zoom in (increase f) and stand further back (increase U) by the same factor

$$D_{TOT} \approx \frac{2NC \boxed{U^2}}{\boxed{f^2}}$$

- ◆ depth of field stays the same, but background gets blurrier!
 - useful for macro when you can't get close enough

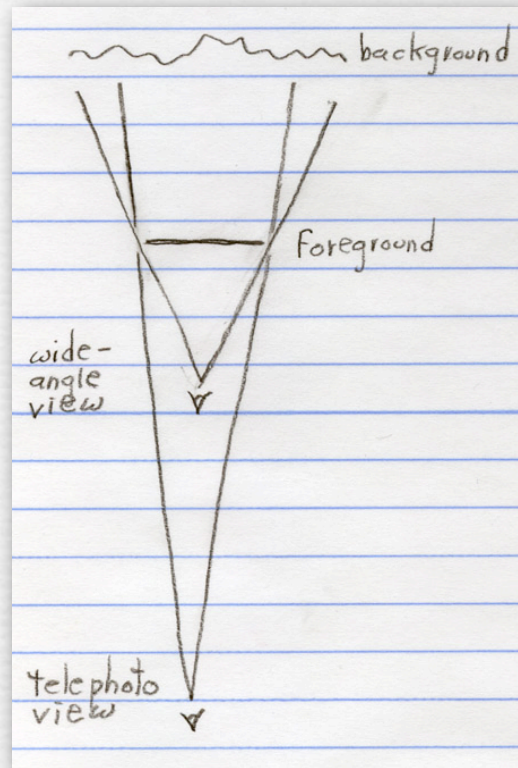


50mm f/4.8



200mm f/4.8,
moved back 4× from subject

Macro photography using a telephoto lens (contents of whiteboard)



- ◆ changing from a wide-angle lens to a telephoto lens and stepping back, you can make a foreground object appear the same size in both lenses
- ◆ and both lenses will have the same depth of field on that object
- ◆ but the telephoto sees a smaller part of the background (which it blows up to fill the field of view), so the background will appear blurrier

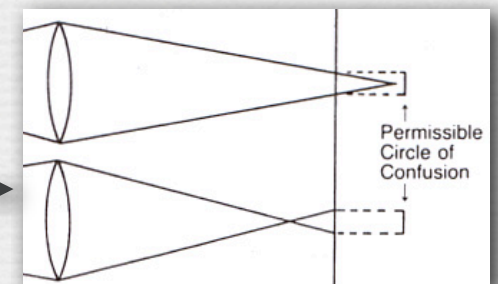
Parting thoughts on DoF: the zen of *bokeh*



Canon 85mm
prime f/1.8 lens



- ◆ the appearance of small out-of-focus features in a photograph with shallow depth of field
 - determined by the boundary of the aperture
 - people get religious about it
 - but not every picture with shallow DoF has evident bokeh...





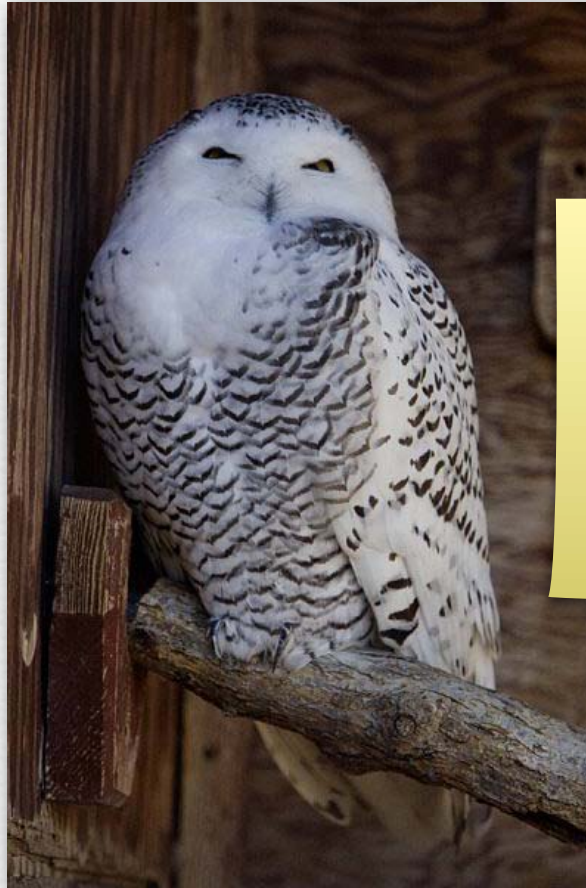
Natasha Gelfand (Canon 100mm f/2.8 prime macro lens)

Games with bokeh



- ◆ picture by Alice Che (CS 178, 2010)
 - heart-shaped mask in front of lens
 - subject was Christmas lights, but misfocused
 - lights were also under-exposed to maintain their color

Parting thoughts on DoF: seeing through occlusions

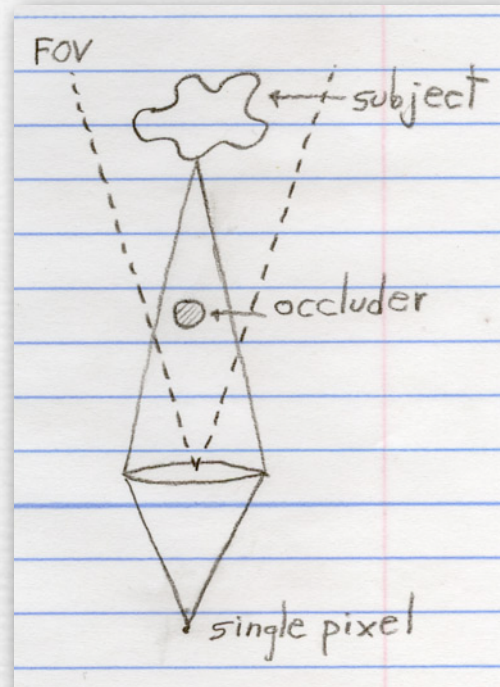


A student asked how seeing through occlusions relates to depth of field. If a camera / lens combination is providing a shallow depth of field, then it is throwing out of focus those scene features that lay beyond the depth of field. Being out of focus makes a scene feature appear blurry, but this is not the same as capturing a sharp photograph of that scene feature, then blurring it in Photoshop. The difference is that optically produced blurring allows you to “see through” occlusions, as diagrammed in the next slide, while blurring producing by convolution (as in Photoshop) does not.

(Fredo Durand)

- ◆ out-of-focus is not the same as convolving the image
 - i.e. not the same as blurring in Photoshop
 - DoF lets you eliminate occlusions, like a chain-link fence

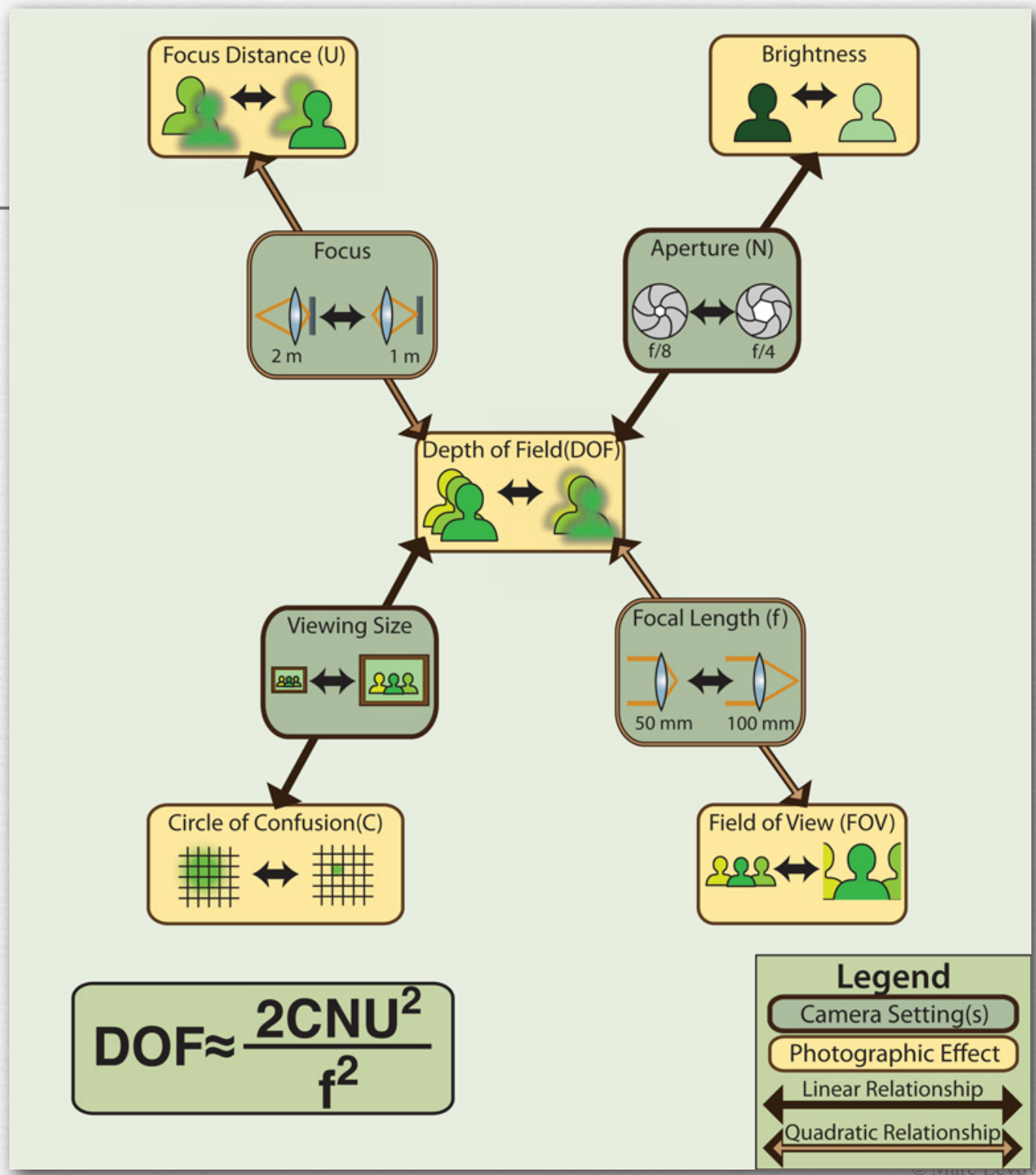
Seeing through occlusions using a large aperture (contents of whiteboard)



- ◆ for a pixel focused on the subject, some of its rays will strike the occluder, but some will pass to the side of it, if the occluder is small enough
- ◆ the pixel will then be a mixture of the colors of the subject and occluder
- ◆ thus, the occluder reduces the contrast of your image of the subject, but it doesn't actually block your view of it

Tradeoffs affecting depth of field

See <http://graphics.stanford.edu/courses/cs178/applets/dof.html>



Recap

- ◆ depth of field (D_{TOT}) is governed by circle of confusion (C), aperture size (N), subject distance (U), and focal length (f)

$$D_{TOT} \approx \frac{2NCU^2}{f^2}$$

- depth of field is linear in some terms and quadratic in others
 - if you focus at the hyperfocal distance $H = f^2 / NC$, everything from $H / 2$ to infinity will be in focus
 - depth of field increases linearly with decreasing sensor size
- ◆ useful sidelights
 - bokeh refers to the appearance of small out-of-focus features
 - you can take macro photographs using a telephoto lens
 - depth of field blur is not the same as blurring an image

Questions?